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Space-time		Line-element	Chapter
Example: Four-dimensional space-time	ds^2	$= -dt^2 + L^2(t,r)dr^2 + B^2(t,r)d\phi^2 + M^2(t,r)dz^2$	7
Hyperbolsk plane – Poincaré Halfplane	ds^2	$= \frac{1}{y^2}(dx^2 + dy^2)$	4,7
Lorentz hyperboloid	ds^2	$= d\psi^2 + \sinh^2 \psi d\theta^2 + \sinh^2 \psi \sin^2 \theta d\phi^2$	7, 13
Poincaré metric	ds^2	$= \frac{4}{1-x^2-y^2}(dx^2 + dy^2)$	7
Rindler metric	ds^2	$= \xi^2 d\tau^2 - d\xi^2$	2,3,4,7
Taub-nut space-time	ds^2	$= -\frac{dt^2}{U^2(t)} + (2l)^2 U^2(t)(dr + \cos \theta d\phi)^2 + V^2(t)(d\theta^2 + \sin^2 \theta d\phi^2)$	7
Three-dimensional flat space in polar coordinates	ds^2	$= dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$	2,7, 13
Tolman-Bondi-de Sitter space-time	ds^2	$= dt^2 - e^{-2\psi(t,r)} dr^2 - R^2(t,r)d\theta^2 - R^2(t,r) \sin^2 \theta d\phi^2$	7
Two-dimensional unit sphere	ds^2	$= d\theta^2 + \sin^2 \theta d\phi^2$	7

7 Cartan's Structure Equations – a Shortcut to the Einstein Tensor

7.1 One-forms.

For general forms, let α be a p -form and β be a q -form, we have

$$\alpha \wedge \beta = (-1)^{pq} \beta \wedge \alpha$$

For one-forms this means

$$\alpha \wedge \beta = -\beta \wedge \alpha$$

$$\Rightarrow \alpha \wedge \alpha = -\alpha \wedge \alpha = 0$$

This also holds for $df = \frac{\partial f}{\partial x^a} dx^a$ because df is a one-form as well.

$$df_a \wedge df_b = \frac{\partial f}{\partial x^a} dx^a \wedge \frac{\partial f}{\partial x^b} dx^b = \frac{\partial^2 f}{\partial x^a \partial x^b} dx^a \wedge dx^b = -\frac{\partial^2 f}{\partial x^a \partial x^b} dx^b \wedge dx^a$$

$$\Rightarrow df_a \wedge df_a = \frac{\partial^2 f}{\partial x^a \partial x^b} dx^a \wedge dx^a$$

and $df_a \wedge df_a = -\frac{\partial^2 f}{\partial x^a \partial x^b} dx^a \wedge dx^a$

Now, because the partial derivatives commute, this can only be true if:

$$dx^a \wedge dx^a = 0$$

7.1.1 The exterior derivative of a one-form.

The exterior derivative of a one-form $f_a dx^a$:

$$d(f_a dx^a) = df_a \wedge dx^a$$

To use this equation it is important to notice, that the right-hand side includes a summation of the partial derivatives times the differential in the usual way:

$$df(x_i) = \frac{\partial}{\partial x_i} f(x_i) dx_i.$$

7.1.2 Examples

$$\sigma = e^{f(r)} dt$$

$$\Rightarrow d\sigma = d(e^{f(r)} dt) = d(e^{f(r)}) \wedge dt = \frac{\partial}{\partial r} (e^{f(r)}) dr \wedge dt = f'(r) e^{f(r)} dr \wedge dt$$

$$\rho = e^{g(r)} \cos \theta \sin \phi dr$$

$$\Rightarrow d\rho = d(e^{g(r)} \cos \theta \sin \phi dr) = d(e^{g(r)} \cos \theta \sin \phi) \wedge dr$$

$$= \frac{\partial}{\partial \theta} (e^{g(r)} \cos \theta \sin \phi) d\theta \wedge dr + \frac{\partial}{\partial \phi} (e^{g(r)} \cos \theta \sin \phi) d\phi \wedge dr$$

$$= -e^{g(r)} \sin \theta \sin \phi d\theta \wedge dr + e^{g(r)} \cos \theta \cos \phi d\phi \wedge dr$$

7.2 Cartan's first structure equation

$$d\omega^{\hat{a}} = -\Gamma^{\hat{a}}_{\hat{b}} \wedge \omega^{\hat{b}}$$

7.3 The curvature two forms and the Riemann tensor

The Riemann tensor and the curvature two forms.

$$\Omega^{\hat{a}}_{\hat{b}} = d\Gamma^{\hat{a}}_{\hat{b}} + \Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} = \frac{1}{2} R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}} \omega^{\hat{c}} \wedge \omega^{\hat{d}}$$

7.4 The unit 2-sphere

The line element:

$$ds^2 = d\theta^2 + \sin^2 \theta d\phi^2$$

7.4.1 The Riemann tensor for the unit 2-sphere

The number of independent elements in the Riemann tensor in a metric of dimension $n = 2$ is $N = \frac{n^2(n^2-1)}{12} = 1$ so we can choose to calculate $R^{\theta}_{\phi\theta\phi}$

The Riemann tensor

$$R^a_{bcd} = \partial_c \Gamma^a_{bd} - \partial_d \Gamma^a_{bc} + \Gamma^e_{bd} \Gamma^a_{ec} - \Gamma^e_{bc} \Gamma^a_{ed}$$

We choose $a = c = \theta$, and $b = d = \phi$

$$R^{\theta}_{\phi\theta\phi} = \partial_{\theta} \Gamma^{\theta}_{\phi\phi} - \partial_{\phi} \Gamma^{\theta}_{\phi\theta} + \Gamma^e_{\phi\phi} \Gamma^{\theta}_{e\theta} - \Gamma^e_{\phi\theta} \Gamma^{\theta}_{e\phi} = \partial_{\theta} \Gamma^{\theta}_{\phi\phi} + \Gamma^e_{\phi\phi} \Gamma^{\theta}_{e\theta} - \Gamma^e_{\phi\theta} \Gamma^{\theta}_{e\phi}$$

$$^1 = \frac{\partial}{\partial r} (e^{g(r)} \cos \theta \sin \phi) dr \wedge dr + \frac{\partial}{\partial \theta} (e^{g(r)} \cos \theta \sin \phi) d\theta \wedge dr + \frac{\partial}{\partial \phi} (e^{g(r)} \cos \theta \sin \phi) d\phi \wedge dr =$$

$$\begin{aligned}
 &= \partial_\theta \Gamma^\theta_{\phi\phi} + \Gamma^\theta_{\phi\phi} \Gamma^\theta_{\theta\theta} - \Gamma^\theta_{\phi\theta} \Gamma^\theta_{\theta\phi} + \Gamma^\phi_{\phi\phi} \Gamma^\theta_{\phi\theta} - \Gamma^\phi_{\phi\theta} \Gamma^\theta_{\phi\phi} = 2\partial_\theta \Gamma^\theta_{\phi\phi} - \Gamma^\phi_{\phi\theta} \Gamma^\theta_{\phi\phi} \\
 &= \partial_\theta (-\sin \theta \cos \theta) - \cot \theta (-\sin \theta \cos \theta) = -\cos^2 \theta + \sin^2 \theta + \cos^2 \theta = \sin^2 \theta
 \end{aligned}$$

7.4.2 The Ricci scalar

The metric tensor and its inverse:

$$\begin{aligned}
 g_{ab} &= \begin{pmatrix} 1 & \\ & \sin^2 \theta \end{pmatrix} \\
 g^{ab} &= \begin{pmatrix} 1 & \\ & \frac{1}{\sin^2 \theta} \end{pmatrix}
 \end{aligned}$$

The Ricci scalar:

$$\begin{aligned}
 R &= g^{ab} R_{ab} = g^{\theta b} R_{\theta b} + g^{\phi b} R_{\phi b} = g^{\theta\theta} R_{\theta\theta} + g^{\theta\phi} R_{\theta\phi} + g^{\phi\theta} R_{\phi\theta} + g^{\phi\phi} R_{\phi\phi} \\
 &= g^{\theta\theta} R_{\theta\theta} + g^{\phi\phi} R_{\phi\phi} = g^{\theta\theta} R^c_{\theta c\theta} + g^{\phi\phi} R^c_{\phi c\phi} \\
 &= g^{\theta\theta} R^\theta_{\theta\theta\theta} + g^{\theta\theta} R^\phi_{\theta\phi\theta} + g^{\phi\phi} R^\theta_{\phi\theta\phi} + g^{\phi\phi} R^\phi_{\phi\phi\phi} = g^{\theta\theta} R^\phi_{\theta\phi\theta} + g^{\phi\phi} R^\theta_{\phi\theta\phi} \\
 &= g^{\theta\theta} g^{\phi\phi} R_{\phi\theta\phi\theta} + g^{\phi\phi} g^{\theta\theta} R_{\theta\phi\theta\phi} = 2g^{\theta\theta} g^{\phi\phi} R_{\theta\phi\theta\phi} = 2g^{\phi\phi} R^\theta_{\phi\theta\phi} = 2 \frac{1}{\sin^2 \theta} \sin^2 \theta = 2
 \end{aligned}$$

7.4.3 Find the Ricci scalar using Cartan's structure equations of the 2-sphere

The line element:

$$ds^2 = d\theta^2 + \sin^2 \theta d\phi^2$$

The Basis one forms:

$$\begin{aligned}
 \omega^{\hat{\theta}} &= d\theta & d\theta &= \omega^{\hat{\theta}} \\
 \omega^{\hat{\phi}} &= \sin \theta d\phi & d\phi &= \frac{1}{\sin \theta} \omega^{\hat{\phi}}
 \end{aligned}$$

$$\eta^{ij} = \begin{pmatrix} 1 & \\ & 1 \end{pmatrix}$$

The curvature one forms:

$$d\omega^{\hat{a}} = -\Gamma^{\hat{a}}_{\hat{b}} \wedge \omega^{\hat{b}}$$

$$d\omega^{\hat{\theta}} = 0$$

$$d\omega^{\hat{\phi}} = d(\sin \theta d\phi) = \cos \theta d\theta \wedge d\phi = \cot \theta \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}}$$

$$\Rightarrow \Gamma^{\hat{\phi}}_{\hat{\theta}} = \cot \theta \omega^{\hat{\phi}}$$

The curvature two forms:

$$\Omega^{\hat{a}}_{\hat{b}} = d\Gamma^{\hat{a}}_{\hat{b}} + \Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} = \frac{1}{2} R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}} \omega^{\hat{c}} \wedge \omega^{\hat{d}}$$

$$d\Gamma^{\hat{\phi}}_{\hat{\theta}} = d(\cot \theta \omega^{\hat{\phi}}) = d(\cos \theta d\phi) = -\sin \theta d\theta \wedge d\phi = -\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}}$$

$$\Gamma^{\hat{\phi}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{\theta}} = \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{\theta}} + \Gamma^{\hat{\phi}}_{\hat{\phi}} \wedge \Gamma^{\hat{\theta}}_{\hat{\theta}} = 0$$

$$\Rightarrow \Omega^{\hat{\phi}}_{\hat{\theta}} = -\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}}$$

$$\Rightarrow R^{\hat{\theta}}_{\hat{\phi}\hat{\theta}\hat{\phi}} = 1$$

$$\Rightarrow R = 2\eta^{\hat{\phi}\hat{\phi}} R^{\hat{\theta}}_{\hat{\phi}\hat{\theta}\hat{\phi}} = 2$$

² We found the Christoffel symbols earlier in the chapter regarding: Two-dimensional sphere with radius a - $ds^2 = a^2 d\theta^2 + a^2 \sin^2 \theta d\phi^2$

³ Notice: $R = 2g^{\phi\phi} R^{\theta}_{\phi\theta\phi}$ is a general solution for a 2-dimensional diagonal metric if we write: $R = 2g^{22} R^1_{212}$

7.5 The three dimensional flat space in spherical polar coordinates

7.5.1 The Christoffel symbols

The line element:

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

The metric tensor and its inverse:

$$g_{ab} = \begin{pmatrix} 1 & & \\ & r^2 & \\ & & r^2 \sin^2 \theta \end{pmatrix}$$

$$g^{ab} = \begin{pmatrix} 1 & & \\ & \frac{1}{r^2} & \\ & & \frac{1}{r^2 \sin^2 \theta} \end{pmatrix}$$

The non-zero Christoffel symbols

$$\begin{aligned} \Gamma_{\theta\theta r} &= -\frac{1}{2} \partial_r g_{\theta\theta} = -\frac{1}{2} \partial_r (r^2) = -r & \Rightarrow \Gamma^r_{\theta\theta} &= g^{rr} \Gamma_{\theta\theta r} = -r \\ \Gamma_{\phi\phi r} &= -\frac{1}{2} \partial_r g_{\phi\phi} = -\frac{1}{2} \partial_r (r^2 \sin^2 \theta) = -r \sin^2 \theta & \Rightarrow \Gamma^r_{\phi\phi} &= g^{rr} \Gamma_{\phi\phi r} = -r \sin^2 \theta \\ \Gamma_{\phi\phi\theta} &= -\frac{1}{2} \partial_\theta g_{\phi\phi} = -r \cos \theta \sin \theta & \Rightarrow \Gamma^\theta_{\phi\phi} &= g^{\theta\theta} \Gamma_{\phi\phi\theta} = -\cos \theta \sin \theta \\ \Gamma_{r\theta\theta} &= \Gamma_{\theta r\theta} = \frac{1}{2} \partial_r g_{\theta\theta} = \frac{1}{2} \partial_r (r^2) = r & \Rightarrow \Gamma^\theta_{r\theta} &= \Gamma^\theta_{\theta r} = g^{\theta\theta} \Gamma_{r\theta\theta} = \frac{1}{r^2} r = \frac{1}{r} \\ \Gamma_{r\phi\phi} &= \Gamma_{\phi r\phi} = \frac{1}{2} \partial_r g_{\phi\phi} = r \sin^2 \theta & \Rightarrow \Gamma^\phi_{r\phi} &= \Gamma^\phi_{\phi r} = g^{\phi\phi} \Gamma_{r\phi\phi} = \frac{1}{r} \\ \Gamma_{\theta\phi\phi} &= \Gamma_{\phi\theta\phi} = \frac{1}{2} \partial_\theta g_{\phi\phi} = r^2 \cos \theta \sin \theta & \Rightarrow \Gamma^\phi_{\theta\phi} &= \Gamma^\phi_{\phi\theta} = g^{\phi\phi} \Gamma_{\theta\phi\phi} = \cot \theta \end{aligned}$$

7.5.2 The Riemann tensor of the three dimensional flat space in spherical polar coordinates

The number of independent elements in the Riemann tensor in a metric of dimension $n = 3$ is $N = \frac{n^2(n^2-1)}{12} = 6$ and we have to calculate: $R^r_{\theta r\theta}$; $R^r_{\phi r\phi}$; $R^\theta_{\phi\theta\phi}$; $R^r_{\theta r\phi}$; $R^\theta_{r\theta\phi}$; $R^\phi_{r\phi\theta}$

The Riemann tensor

$$\begin{aligned} R^a_{bcd} &= \partial_c \Gamma^a_{bd} - \partial_d \Gamma^a_{bc} + \Gamma^e_{bd} \Gamma^a_{ec} - \Gamma^e_{bc} \Gamma^a_{ed} \\ R^r_{\theta r\theta} &= \partial_r \Gamma^r_{\theta\theta} - \partial_\theta \Gamma^r_{\theta r} + \Gamma^e_{\theta\theta} \Gamma^r_{er} - \Gamma^e_{\theta r} \Gamma^r_{e\theta} = \partial_r \Gamma^r_{\theta\theta} - \Gamma^\theta_{\theta r} \Gamma^r_{\theta\theta} = -1 - \left(\frac{1}{r}\right)(-r) \\ &= 0 \\ R^r_{\phi r\phi} &= \partial_r \Gamma^r_{\phi\phi} - \partial_\phi \Gamma^r_{\phi r} + \Gamma^e_{\phi\phi} \Gamma^r_{er} - \Gamma^e_{\phi r} \Gamma^r_{e\phi} = \partial_r \Gamma^r_{\phi\phi} - \Gamma^\phi_{\phi r} \Gamma^r_{\phi\phi} \\ &= -\sin^2 \theta - \left(\frac{1}{r}\right)(-r \sin^2 \theta) = 0 \\ R^\theta_{\phi\theta\phi} &= \partial_\theta \Gamma^\theta_{\phi\phi} - \partial_\phi \Gamma^\theta_{\phi\theta} + \Gamma^e_{\phi\phi} \Gamma^\theta_{e\theta} - \Gamma^e_{\phi\theta} \Gamma^\theta_{e\phi} = \partial_\theta \Gamma^\theta_{\phi\phi} + \Gamma^r_{\phi\phi} \Gamma^\theta_{r\theta} - \Gamma^\phi_{\phi\theta} \Gamma^\theta_{\phi\phi} \\ &= -\cos^2 \theta + \sin^2 \theta + (-r \sin^2 \theta) \left(\frac{1}{r}\right) - (\cot \theta)(-\sin \theta \cos \theta) = 0 \end{aligned}$$

$$^4 = -\frac{1}{2} \partial_r (r^2 \sin^2 \theta) =$$

$$^5 = \frac{1}{r^2} (-r^2 \cos \theta \sin \theta) =$$

$$^6 = \frac{1}{2} \partial_r (r^2 \sin^2 \theta) =$$

$$^7 = \frac{1}{r^2 \sin^2 \theta} r \sin^2 \theta =$$

$$^8 = \frac{1}{2} \partial_\theta (r^2 \sin^2 \theta) =$$

$$^9 = \frac{1}{r^2 \sin^2 \theta} r^2 \cos \theta \sin \theta =$$

$$\begin{aligned}
 R^r_{\theta r \phi} &= \partial_r \Gamma^r_{\theta \phi} - \partial_\phi \Gamma^r_{\theta r} + \Gamma^e_{\theta \phi} \Gamma^r_{er} - \Gamma^e_{\theta r} \Gamma^r_{e\phi} = 0 \\
 R^\theta_{r \theta \phi} &= \partial_\theta \Gamma^\theta_{r \phi} - \partial_\phi \Gamma^\theta_{r \theta} + \Gamma^e_{r \phi} \Gamma^\theta_{e\theta} - \Gamma^e_{r \theta} \Gamma^\theta_{e\phi} = 0 \\
 R^\phi_{r \phi \theta} &= \partial_\phi \Gamma^\phi_{r \theta} - \partial_\theta \Gamma^\phi_{r \phi} + \Gamma^e_{r \theta} \Gamma^\phi_{e\phi} - \Gamma^e_{r \phi} \Gamma^\phi_{e\theta} = \left(\frac{1}{r}\right) \cot \theta - \left(\frac{1}{r}\right) \cot \theta = 0
 \end{aligned}$$

Not surprisingly all the elements of the Riemann tensor in flat three-dimensional space are zero.

7.5.3 The Ricci rotation coefficients

The line element:

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

The Basis one forms and the transformation matrices

$$\begin{aligned}
 \omega^{\hat{r}} &= dr & dr &= \omega^{\hat{r}} \\
 \omega^{\hat{\theta}} &= r d\theta & d\theta &= \frac{1}{r} \omega^{\hat{\theta}} \\
 \omega^{\hat{\phi}} &= r \sin \theta d\phi & d\phi &= \frac{1}{r \sin \theta} \omega^{\hat{\phi}}
 \end{aligned}$$

$$\begin{aligned}
 \eta^{ij} &= \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix} \\
 \Lambda^{\hat{a}}_b &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & r & 0 \\ 0 & 0 & r \sin \theta \end{pmatrix} \\
 \Lambda^a_{\hat{b}} &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{1}{r} & 0 \\ 0 & 0 & \frac{1}{r \sin \theta} \end{pmatrix}
 \end{aligned}$$

The Ricci rotation coefficients $\Gamma^{\hat{a}}_{\hat{b}\hat{c}}$

$$\begin{aligned}
 d\omega^{\hat{a}} &= -\Gamma^{\hat{a}}_{\hat{b}\hat{c}} \omega^{\hat{b}} \omega^{\hat{c}} \\
 \Gamma^{\hat{a}}_{\hat{b}\hat{c}} &= \Gamma^{\hat{a}}_{\hat{b}\hat{c}} \omega^{\hat{c}} \\
 d\omega^{\hat{r}} &= 0 \\
 d\omega^{\hat{\theta}} &= d(r d\theta) = dr \wedge d\theta = \frac{1}{r} \omega^{\hat{r}} \wedge \omega^{\hat{\theta}} = -\frac{1}{r} \omega^{\hat{\theta}} \wedge \omega^{\hat{r}} = -\Gamma^{\hat{\theta}}_{\hat{r}\hat{\theta}} \omega^{\hat{r}} \omega^{\hat{\theta}} \\
 d\omega^{\hat{\phi}} &= d(r \sin \theta d\phi) = \sin \theta dr \wedge d\phi + r \cos \theta d\theta \wedge d\phi = \frac{1}{r} \omega^{\hat{r}} \wedge \omega^{\hat{\phi}} + \frac{\cot \theta}{r} \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\
 &= -\frac{1}{r} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} - \frac{\cot \theta}{r} \omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} = -\Gamma^{\hat{\phi}}_{\hat{r}\hat{\phi}} \omega^{\hat{r}} \omega^{\hat{\phi}} - \Gamma^{\hat{\phi}}_{\hat{\theta}\hat{\phi}} \omega^{\hat{\theta}} \omega^{\hat{\phi}}
 \end{aligned}$$

Summarizing the curvature one forms in a matrix (where \hat{a} refers to column and \hat{b} to row):

$$\Gamma^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 0 & \frac{1}{r} \omega^{\hat{\theta}} & \frac{1}{r} \omega^{\hat{\phi}} \\ -\frac{1}{r} \omega^{\hat{\theta}} & 0 & \frac{\cot \theta}{r} \omega^{\hat{\phi}} \\ -\frac{1}{r} \omega^{\hat{\phi}} & -\frac{\cot \theta}{r} \omega^{\hat{\phi}} & 0 \end{pmatrix}$$

Now we can read off the Ricci rotation coefficients

$$\begin{aligned}
 \Gamma^{\hat{r}}_{\hat{\theta}\hat{\theta}} &= -\frac{1}{r} & \Gamma^{\hat{\theta}}_{\hat{r}\hat{\theta}} &= \frac{1}{r} & \Gamma^{\hat{\phi}}_{\hat{r}\hat{\phi}} &= \frac{1}{r} \\
 \Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}} &= -\frac{1}{r} & \Gamma^{\hat{\theta}}_{\hat{\phi}\hat{\phi}} &= -\frac{\cot \theta}{r} & \Gamma^{\hat{\phi}}_{\hat{\theta}\hat{\phi}} &= \frac{\cot \theta}{r}
 \end{aligned}$$

7.5.4 Transformation of the Ricci rotation coefficients $\Gamma^{\hat{a}}_{\hat{b}\hat{c}}$ into the Christoffel symbols Γ^a_{bc} of the three dimensional flat space in spherical polar coordinates

The transformation

$$\begin{aligned} \Gamma^a_{bc} &= \Lambda^a_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_b \Lambda^{\hat{f}}_c \\ \Rightarrow \Gamma^{\theta}_{r\theta} &= \Lambda^{\theta}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_r \Lambda^{\hat{f}}_{\theta} = \Gamma^{\hat{\theta}}_{\hat{r}\hat{\theta}} \Lambda^{\hat{r}}_r = \frac{1}{r} \cdot 1 = \frac{1}{r} \\ \Gamma^r_{\theta\theta} &= \Lambda^r_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\theta} \Lambda^{\hat{f}}_{\theta} = \Lambda^r_{\hat{r}} \Gamma^{\hat{r}}_{\hat{\theta}\hat{\theta}} (\Lambda^{\hat{\theta}}_{\theta})^2 = 1 \left(-\frac{1}{r}\right) r^2 = -r \\ \Gamma^{\phi}_{r\phi} &= \Lambda^{\phi}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_r \Lambda^{\hat{f}}_{\phi} = \Gamma^{\hat{\phi}}_{\hat{r}\hat{\phi}} \Lambda^{\hat{r}}_r = \frac{1}{r} \cdot 1 = \frac{1}{r} \\ \Gamma^r_{\phi\phi} &= \Lambda^r_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\phi} \Lambda^{\hat{f}}_{\phi} = \Lambda^r_{\hat{r}} \Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}} (\Lambda^{\hat{\phi}}_{\phi})^2 = 1 \left(-\frac{1}{r}\right) r^2 \sin^2 \theta = -r \sin^2 \theta \\ \Gamma^{\phi}_{\theta\phi} &= \Lambda^{\phi}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\theta} \Lambda^{\hat{f}}_{\phi} = \Gamma^{\hat{\phi}}_{\hat{\theta}\hat{\phi}} \Lambda^{\hat{\theta}}_{\theta} = \frac{\cot \theta}{r} \cdot r = \cot \theta \\ \Gamma^{\theta}_{\phi\phi} &= \Lambda^{\theta}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\phi} \Lambda^{\hat{f}}_{\phi} = \Lambda^{\theta}_{\hat{\theta}} \Gamma^{\hat{\theta}}_{\hat{\phi}\hat{\phi}} (\Lambda^{\hat{\phi}}_{\phi})^2 = \frac{1}{r} \left(-\frac{\cot \theta}{r}\right) r^2 \sin^2 \theta = -\sin \theta \cos \theta \end{aligned}$$

Collecting the results

$$\begin{aligned} \Gamma^r_{\theta\theta} &= -r & \Gamma^r_{\phi\phi} &= -r \sin^2 \theta & \Gamma^{\theta}_{\phi\phi} &= -\sin \theta \cos \theta \\ \Gamma^{\theta}_{r\theta} &= \frac{1}{r} & \Gamma^{\phi}_{r\phi} &= \frac{1}{r} & \Gamma^{\phi}_{\theta\phi} &= \cot \theta \end{aligned}$$

Which is consistent with the ordinary method used above.

7.6 The Rindler metric

¹⁰The Rindler coordinate system or frame describes a uniformly accelerating frame of reference in Minkowski space.

7.6.1 Use the geodesic equations to find the Christoffel symbols for the Rindler metric.

The line element:

$$ds^2 = \xi^2 d\tau^2 - d\xi^2$$

The metric tensor:

$$g_{ab} = \begin{Bmatrix} \xi^2 & \\ & -1 \end{Bmatrix}$$

The geodesic equation

$$\begin{aligned} \frac{\partial K}{\partial x^a} &= \frac{d}{ds} \left(\frac{\partial K}{\partial \dot{x}^a} \right) \\ K &= \frac{1}{2} g_{ab} \dot{x}^a \dot{x}^b = \frac{1}{2} \xi^2 (\dot{\tau})^2 - \frac{1}{2} \dot{\xi}^2 \end{aligned}$$

$\underline{x^a \equiv \xi:}$ $\frac{\partial K}{\partial \dot{\xi}} = \frac{d}{ds} \left(\frac{\partial K}{\partial \dot{\xi}} \right)$

$$\Rightarrow \xi \dot{\tau}^2 = \frac{d}{ds} (-\dot{\xi}) = -\ddot{\xi}$$

$$\Rightarrow 0 = \ddot{\xi} + \xi \dot{\tau}^2$$

$\underline{x^a \equiv \tau:}$ $\frac{\partial K}{\partial \dot{\tau}} = \frac{d}{ds} \left(\frac{\partial K}{\partial \dot{\tau}} \right)$

$$0 = \frac{d}{ds} (\xi^2 \dot{\tau}) = 2\xi \dot{\xi} \dot{\tau} + \xi^2 \ddot{\tau}$$

$$\Rightarrow 0 = \ddot{\tau} + \frac{1}{\xi} \dot{\xi} \dot{\tau} + \frac{1}{\xi} \dot{\tau} \dot{\xi}$$

¹⁰ We have looked at this space-time earlier: $ds^2 = -X^2 dT^2 + dX^2$

Collecting the results

$$\begin{aligned} 0 &= \ddot{\xi} + \xi \dot{t}^2 \\ 0 &= 2\xi \dot{\xi} \dot{t} + \xi^2 \ddot{t} \\ \Rightarrow \Gamma_{\tau\tau}^{\xi} &= \xi \\ \Gamma_{\xi\tau}^{\tau} &= \frac{1}{\xi} \\ \Gamma_{\tau\xi}^{\tau} &= \frac{1}{\xi} \end{aligned}$$

7.6.2 Ricci rotation coefficients of the Rindler metric

The line element:

$$ds^2 = u^2 dv^2 - du^2$$

The Basis one forms

$$\begin{aligned} \omega^{\hat{u}} &= du & du &= \omega^{\hat{u}} \\ \omega^{\hat{v}} &= u dv & dv &= \frac{1}{u} \omega^{\hat{v}} \end{aligned}$$

$$\eta^{ij} = \begin{Bmatrix} 1 & \\ & -1 \end{Bmatrix}$$

Cartan's First Structure equation and the calculation of the Ricci rotation coefficients $\Gamma_{\hat{b}\hat{c}}^{\hat{a}}$:

$$\begin{aligned} d\omega^{\hat{a}} &= -\Gamma_{\hat{b}\hat{c}}^{\hat{a}} \omega^{\hat{b}} \wedge \omega^{\hat{c}} \\ \Rightarrow d\omega^{\hat{u}} &= 0 \\ d\omega^{\hat{v}} &= d(u dv) = du \wedge dv = \frac{1}{u} \omega^{\hat{u}} \wedge \omega^{\hat{v}} \\ \Rightarrow \Gamma_{\hat{u}\hat{u}}^{\hat{v}} &= \frac{1}{u} \\ \Gamma_{\hat{u}\hat{v}}^{\hat{v}} &= \frac{1}{u} \\ \Gamma_{\hat{v}\hat{v}}^{\hat{u}} &= \eta^{\hat{u}\hat{u}} \Gamma_{\hat{u}\hat{v}\hat{v}} = -\eta^{\hat{u}\hat{u}} \Gamma_{\hat{v}\hat{u}\hat{v}} = -\eta^{\hat{u}\hat{u}} \eta_{\hat{v}\hat{v}} \Gamma_{\hat{u}\hat{v}}^{\hat{v}} = \frac{1}{u} \end{aligned}$$

7.7 The Lorentz hyperboloid

7.7.1 The Christoffel symbols.

The line element:

$$ds^2 = d\psi^2 + \sinh^2 \psi d\theta^2 + \sinh^2 \psi \sin^2 \theta d\phi^2$$

To find the Christoffel symbols we calculate the geodesic from the Euler-Lagrange equation

$$\begin{aligned} 0 &= \frac{d}{ds} \left(\frac{\partial F}{\partial \dot{x}^a} \right) - \frac{\partial F}{\partial x^a} \\ \Rightarrow F &= \dot{\psi}^2 + \sinh^2 \psi \dot{\theta}^2 + \sinh^2 \psi \sin^2 \theta \dot{\phi}^2 \\ \underline{x^a = \psi}: \quad \frac{\partial F}{\partial \dot{\psi}} &= 2 \cosh \psi \sinh \psi \dot{\theta}^2 + 2 \cosh \psi \sinh \psi \sin^2 \theta \dot{\phi}^2 \\ \frac{\partial F}{\partial \dot{\psi}} &= 2\dot{\psi} \\ \frac{d}{ds} \left(\frac{\partial F}{\partial \dot{\psi}} \right) &= 2\ddot{\psi} \\ \Rightarrow 0 &= 2\ddot{\psi} - 2 \cosh \psi \sinh \psi \dot{\theta}^2 - 2 \cosh \psi \sinh \psi \sin^2 \theta \dot{\phi}^2 \\ \Rightarrow 0 &= \ddot{\psi} - \cosh \psi \sinh \psi \dot{\theta}^2 - \cosh \psi \sinh \psi \sin^2 \theta \dot{\phi}^2 \\ \underline{x^a = \theta}: \quad \frac{\partial F}{\partial \dot{\theta}} &= 2 \cos \theta \sin \theta \sinh^2 \psi \dot{\phi}^2 \end{aligned}$$

$$\begin{aligned} \frac{\partial F}{\partial \theta} &= 2 \sinh^2 \psi \dot{\theta} \\ \frac{d}{ds} \left(\frac{\partial F}{\partial \dot{\theta}} \right) &= 4 \cosh \psi \sinh \psi \dot{\psi} \dot{\theta} + 2 \sinh^2 \psi \ddot{\theta} \\ \Rightarrow 0 &= 2 \sinh^2 \psi \ddot{\theta} + 4 \cosh \psi \sinh \psi \dot{\psi} \dot{\theta} - 2 \cos \theta \sin \theta \sinh^2 \psi \dot{\phi}^2 \\ \Rightarrow 0 &= \ddot{\theta} + 2 \coth \psi \dot{\psi} \dot{\theta} - \cos \theta \sin \theta \dot{\phi}^2 \\ \underline{x^a = \phi}: \quad \frac{\partial F}{\partial \dot{\phi}} &= 0 \\ \frac{\partial F}{\partial \dot{\phi}} &= 2 \sinh^2 \psi \sin^2 \theta \dot{\phi} \\ \frac{d}{ds} \left(\frac{\partial F}{\partial \dot{\phi}} \right) &= 4 \cosh \psi \sinh \psi \dot{\psi} \dot{\phi} + 4 \cos \theta \sin \theta \dot{\theta} \dot{\phi} + 2 \sinh^2 \psi \sin^2 \theta \ddot{\phi} \\ \Rightarrow 0 &= 4 \cosh \psi \sinh \psi \dot{\psi} \dot{\phi} + 4 \cos \theta \sin \theta \dot{\theta} \dot{\phi} + 2 \sinh^2 \psi \sin^2 \theta \ddot{\phi} \\ \Rightarrow 0 &= \ddot{\phi} + 2 \frac{\coth \psi}{\sin^2 \theta} \dot{\psi} \dot{\phi} + 2 \frac{\cot \theta}{\sinh^2 \psi} \dot{\theta} \dot{\phi} \end{aligned}$$

Collecting the results

$$\begin{aligned} 0 &= \ddot{\psi} - \cosh \psi \sinh \psi \dot{\theta}^2 - \cosh \psi \sinh \psi \sin^2 \theta \dot{\phi}^2 \\ 0 &= \ddot{\theta} + 2 \cosh \psi \sinh \psi \dot{\psi} \dot{\theta} - \cos \theta \sin \theta \dot{\phi}^2 \\ 0 &= \ddot{\phi} + 2 \frac{\coth \psi}{\sin^2 \theta} \dot{\psi} \dot{\phi} + 2 \frac{\cot \theta}{\sinh^2 \psi} \dot{\theta} \dot{\phi} \end{aligned}$$

We can now find the non-zero Christoffel symbols:

$$\begin{aligned} \Gamma_{\theta\theta}^{\psi} &= -\cosh \psi \sinh \psi \\ \Gamma_{\phi\phi}^{\psi} &= -\cosh \psi \sinh \psi \sin^2 \theta \\ \Gamma_{\psi\theta}^{\theta} &= \Gamma_{\theta\psi}^{\theta} = \cosh \psi \sinh \psi \\ \Gamma_{\phi\phi}^{\theta} &= -\cos \theta \sin \theta \\ \Gamma_{\psi\phi}^{\phi} &= \Gamma_{\phi\psi}^{\phi} = \frac{\coth \psi}{\sin^2 \theta} \\ \Gamma_{\theta\phi}^{\phi} &= \Gamma_{\phi\theta}^{\phi} = \frac{\cot \theta}{\sinh^2 \psi} \end{aligned}$$

7.7.2 The Ricci rotation coefficients

The line element:

$$ds^2 = d\psi^2 + \sinh^2 \psi d\theta^2 + \sinh^2 \psi \sin^2 \theta d\phi^2$$

The Basis one forms

$$\begin{aligned} \omega^{\hat{\psi}} &= d\psi & d\psi &= \omega^{\hat{\psi}} \\ \omega^{\hat{\theta}} &= \sinh \psi d\theta & d\theta &= \frac{1}{\sinh \psi} \omega^{\hat{\theta}} \\ \omega^{\hat{\phi}} &= \sinh \psi \sin \theta d\phi & d\phi &= \frac{1}{\sinh \psi \sin \theta} \omega^{\hat{\phi}} \end{aligned}$$

$$\eta^{ij} = \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix}$$

Cartan's First Structure equation and the calculation of the Ricci rotation coefficients $\Gamma_{\hat{b}\hat{c}}^{\hat{a}}$:

$$\begin{aligned} d\omega^{\hat{a}} &= -\Gamma_{\hat{b}\hat{c}}^{\hat{a}} \omega^{\hat{b}} \wedge \omega^{\hat{c}} \\ \Gamma_{\hat{b}\hat{c}}^{\hat{a}} &= \Gamma_{\hat{b}\hat{c}}^{\hat{a}} \omega^{\hat{c}} \\ \Rightarrow d\omega^{\hat{\psi}} &= 0 \\ d\omega^{\hat{\theta}} &= d(\sinh \psi d\theta) = \cosh \psi d\psi \wedge d\theta = \cosh \psi \omega^{\hat{\psi}} \wedge \frac{1}{\sinh \psi} \omega^{\hat{\theta}} = \coth \psi \omega^{\hat{\psi}} \wedge \omega^{\hat{\theta}} \end{aligned}$$

$$\begin{aligned}
 d\omega^{\hat{\phi}} &= d(\sinh \psi \sin \theta d\phi) = \cosh \psi d\psi \wedge d\phi + \cos \theta d\theta \wedge d\phi \\
 &= \cosh \psi \omega^{\hat{\psi}} \wedge \frac{1}{\sinh \psi \sin \theta} \omega^{\hat{\phi}} + \cos \theta \frac{1}{\sinh \psi} \omega^{\hat{\theta}} \wedge \frac{1}{\sinh \psi \sin \theta} \omega^{\hat{\phi}} \\
 &= \frac{\coth \psi}{\sin \theta} \omega^{\hat{\psi}} \wedge \omega^{\hat{\phi}} + \frac{\cot \theta}{\sinh^2 \psi} \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}}
 \end{aligned}$$

Summarizing the curvature one forms in a matrix:

$$\Gamma^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 0 & \coth \psi \omega^{\hat{\theta}} & \frac{\coth \psi}{\sin \theta} \omega^{\hat{\phi}} \\ -\coth \psi \omega^{\hat{\theta}} & 0 & \frac{\cot \theta}{\sinh^2 \psi} \omega^{\hat{\phi}} \\ -\frac{\coth \psi}{\sin \theta} \omega^{\hat{\phi}} & -\frac{\cot \theta}{\sinh^2 \psi} \omega^{\hat{\phi}} & 0 \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row.

Now we can find the non-zero Ricci rotation coefficients:

$$\begin{aligned}
 \Gamma^{\hat{\psi}}_{\hat{\theta}\hat{\theta}} &= -\coth \psi \\
 \Gamma^{\hat{\psi}}_{\hat{\phi}\hat{\phi}} &= -\frac{\coth \psi}{\sin \theta} \\
 \Gamma^{\hat{\theta}}_{\hat{\psi}\hat{\theta}} &= \coth \psi \\
 \Gamma^{\hat{\theta}}_{\hat{\phi}\hat{\phi}} &= -\frac{\cot \theta}{\sinh^2 \psi} \\
 \Gamma^{\hat{\phi}}_{\hat{\psi}\hat{\phi}} &= \frac{\coth \psi}{\sin \theta} \\
 \Gamma^{\hat{\phi}}_{\hat{\theta}\hat{\phi}} &= \frac{\cot \theta}{\sinh^2 \psi}
 \end{aligned}$$

7.8 Ricci rotation coefficients, Ricci scalar and Einstein equations for a general 4-dimensional metric: $ds^2 = -dt^2 + L^2(t,r)dr^2 + B^2(t,r)d\phi^2 + M^2(t,r)dz^2$

The line element:

$$ds^2 = -dt^2 + L^2(t,r)dr^2 + B^2(t,r)d\phi^2 + M^2(t,r)dz^2$$

7.8.1 The Basis one forms

$$\begin{aligned}
 \omega^{\hat{t}} &= dt \\
 \omega^{\hat{r}} &= L(t,r)dr & dr &= \frac{1}{L(t,r)}\omega^{\hat{r}} \\
 \omega^{\hat{\phi}} &= B(t,r)d\phi & d\phi &= \frac{1}{B(t,r)}\omega^{\hat{\phi}} \\
 \omega^{\hat{z}} &= M(t,r)dz & dz &= \frac{1}{M(t,r)}\omega^{\hat{z}} \\
 \eta^{ij} &= \begin{pmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{pmatrix}
 \end{aligned}$$

7.8.2 Cartan's First Structure equation and the calculation of the Ricci rotation coefficients $\Gamma^{\hat{a}}_{\hat{b}\hat{c}}$

$$\begin{aligned}
 d\omega^{\hat{a}} &= -\Gamma^{\hat{a}}_{\hat{b}\hat{c}} \omega^{\hat{b}} \wedge \omega^{\hat{c}} \\
 \Gamma^{\hat{a}}_{\hat{b}\hat{c}} &= \Gamma^{\hat{a}}_{\hat{c}\hat{b}} \\
 d\omega^{\hat{t}} &= 0
 \end{aligned}$$

$$d\omega^{\hat{r}} = d(L(t,r)dr) = \dot{L}dt \wedge dr = \frac{\dot{L}}{L}\omega^{\hat{t}} \wedge \omega^{\hat{r}}$$

$$d\omega^{\hat{\phi}} = d(B(t,r)d\phi) = \dot{B}dt \wedge d\phi + B'dr \wedge d\phi = \frac{\dot{B}}{B}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{B'}{LB}\omega^{\hat{r}} \wedge \omega^{\hat{\phi}}$$

$$d\omega^{\hat{z}} = d(M(t,r)dz) = \dot{M}dt \wedge dz + M'dr \wedge dz = \frac{\dot{M}}{M}\omega^{\hat{t}} \wedge \omega^{\hat{z}} + \frac{M'}{LM}\omega^{\hat{r}} \wedge \omega^{\hat{z}}$$

Summarizing the curvature one forms in a matrix:

$$\Gamma^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 0 & \frac{\dot{L}}{L}\omega^{\hat{r}} & \frac{\dot{B}}{B}\omega^{\hat{\phi}} & \frac{\dot{M}}{M}\omega^{\hat{z}} \\ \frac{\dot{L}}{L}\omega^{\hat{r}} & 0 & \frac{B'}{LB}\omega^{\hat{\phi}} & \frac{M'}{LM}\omega^{\hat{z}} \\ \frac{\dot{B}}{B}\omega^{\hat{\phi}} & -\frac{B'}{LB}\omega^{\hat{\phi}} & 0 & 0 \\ \frac{\dot{M}}{M}\omega^{\hat{z}} & -\frac{M'}{LM}\omega^{\hat{z}} & 0 & 0 \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row

Now we can read off the Ricci rotation coefficients

$$\begin{aligned} \Gamma^{\hat{t}}_{\hat{r}\hat{r}} &= \frac{\dot{L}}{L} & \Gamma^{\hat{r}}_{\hat{t}\hat{r}} &= \frac{\dot{L}}{L} & \Gamma^{\hat{\phi}}_{\hat{t}\hat{\phi}} &= \frac{\dot{B}}{B} & \Gamma^{\hat{z}}_{\hat{t}\hat{z}} &= \frac{\dot{M}}{M} \\ \Gamma^{\hat{t}}_{\hat{\phi}\hat{\phi}} &= \frac{\dot{B}}{B} & \Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}} &= -\frac{B'}{LB} & \Gamma^{\hat{\phi}}_{\hat{r}\hat{\phi}} &= \frac{B'}{LB} & \Gamma^{\hat{z}}_{\hat{r}\hat{z}} &= \frac{M'}{LM} \\ \Gamma^{\hat{t}}_{\hat{z}\hat{z}} &= \frac{\dot{M}}{M} & \Gamma^{\hat{r}}_{\hat{z}\hat{z}} &= -\frac{M'}{LM} \end{aligned}$$

7.8.3 The curvature two forms

$$\Omega^{\hat{a}}_{\hat{b}} = d\Gamma^{\hat{a}}_{\hat{b}} + \Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} = \frac{1}{2}R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}}\omega^{\hat{c}} \wedge \omega^{\hat{d}}$$

$$d\Gamma^{\hat{r}}_{\hat{t}} = d\left(\frac{\dot{L}}{L}\omega^{\hat{r}}\right) = d(\dot{L}(t,r)dr) = \ddot{L}dt \wedge dr + \dot{L}'dr \wedge dr = \frac{\ddot{L}}{L}\omega^{\hat{t}} \wedge \omega^{\hat{r}}$$

$$\Gamma^{\hat{r}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} = \Gamma^{\hat{r}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{z}} \wedge \Gamma^{\hat{z}}_{\hat{t}} = 0$$

$$\Rightarrow \Omega^{\hat{r}}_{\hat{t}} = -\frac{\ddot{L}}{L}\omega^{\hat{r}} \wedge \omega^{\hat{t}}$$

$$d\Gamma^{\hat{\phi}}_{\hat{t}} = d\left(\frac{\dot{B}}{B}\omega^{\hat{\phi}}\right) = d(\dot{B}(t,r)d\phi) = \ddot{B}dt \wedge d\phi + \dot{B}'dr \wedge d\phi$$

$$= \frac{\ddot{B}}{B}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{\dot{B}'}{BL}\omega^{\hat{r}} \wedge \omega^{\hat{\phi}}$$

$$\Gamma^{\hat{\phi}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} = \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{z}} \wedge \Gamma^{\hat{z}}_{\hat{t}} = \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} = \frac{B'}{BL}\omega^{\hat{\phi}} \wedge \frac{\dot{L}}{L}\omega^{\hat{r}}$$

$$\Rightarrow \Omega^{\hat{\phi}}_{\hat{t}} = \frac{\ddot{B}}{B}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{\dot{B}'}{BL}\omega^{\hat{r}} \wedge \omega^{\hat{\phi}} + \frac{B'}{BL}\omega^{\hat{\phi}} \wedge \frac{\dot{L}}{L}\omega^{\hat{r}}$$

$$= -\frac{\ddot{B}}{B}\omega^{\hat{\phi}} \wedge \omega^{\hat{t}} + \left(\frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL}\right)\omega^{\hat{\phi}} \wedge \omega^{\hat{r}}$$

$$d\Gamma^{\hat{z}}_{\hat{t}} = d\left(\frac{\dot{M}}{M}\omega^{\hat{z}}\right) = d(\dot{M}(t,r)dz) = \ddot{M}dt \wedge dz + \dot{M}'dr \wedge dz$$

$$= \frac{\ddot{M}}{M}\omega^{\hat{t}} \wedge \omega^{\hat{z}} + \frac{\dot{M}'}{ML}\omega^{\hat{r}} \wedge \omega^{\hat{z}}$$

$$\begin{aligned}
 \Gamma^{\hat{z}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} &= \Gamma^{\hat{z}}_{\hat{t}} \wedge \Gamma^{\hat{c}}_{\hat{t}} + \Gamma^{\hat{z}}_{\hat{r}} \wedge \Gamma^{\hat{c}}_{\hat{t}} + \Gamma^{\hat{z}}_{\hat{\phi}} \wedge \Gamma^{\hat{c}}_{\hat{t}} + \Gamma^{\hat{z}}_{\hat{z}} \wedge \Gamma^{\hat{c}}_{\hat{t}} = \Gamma^{\hat{z}}_{\hat{r}} \wedge \Gamma^{\hat{c}}_{\hat{t}} = \frac{M'}{ML} \omega^{\hat{z}} \wedge \frac{\dot{L}}{L} \omega^{\hat{r}} \\
 \Rightarrow \Omega^{\hat{z}}_{\hat{t}} &= \frac{\ddot{M}}{M} \omega^{\hat{t}} \wedge \omega^{\hat{z}} + \frac{\dot{M}'}{ML} \omega^{\hat{r}} \wedge \omega^{\hat{z}} + \frac{M'}{ML} \omega^{\hat{z}} \wedge \frac{\dot{L}}{L} \omega^{\hat{r}} \\
 &= -\frac{\ddot{M}}{M} \omega^{\hat{z}} \wedge \omega^{\hat{t}} + \left(\frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} \right) \omega^{\hat{z}} \wedge \omega^{\hat{r}} \\
 d\Gamma^{\hat{\phi}}_{\hat{r}} &= d\left(\frac{B'}{BL} \omega^{\hat{\phi}} \right) = d\left(\frac{B'}{L} d\phi \right) = \frac{\dot{B}'L - B'\dot{L}}{L^2} dt \wedge d\phi + \frac{B''L - B'L'}{L^2} dr \wedge d\phi \\
 &= \frac{B'\dot{L} - \dot{B}'L}{BL^2} \omega^{\hat{\phi}} \wedge \omega^{\hat{t}} + \frac{B'L' - B''L}{BL^3} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\
 \Gamma^{\hat{\phi}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{r}} &= \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{c}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{c}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{\phi}} \wedge \Gamma^{\hat{c}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{z}} \wedge \Gamma^{\hat{c}}_{\hat{r}} = \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{c}}_{\hat{r}} = \frac{\dot{B}\dot{L}}{BL} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\
 \Rightarrow \Omega^{\hat{\phi}}_{\hat{r}} &= \frac{B'\dot{L} - \dot{B}'L}{BL^2} \omega^{\hat{\phi}} \wedge \omega^{\hat{t}} + \frac{B'L' - B''L}{BL^3} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} + \frac{\dot{B}\dot{L}}{BL} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\
 &= \frac{B'\dot{L} - \dot{B}'L}{BL^2} \omega^{\hat{\phi}} \wedge \omega^{\hat{t}} + \left(\frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} \right) \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\
 d\Gamma^{\hat{z}}_{\hat{r}} &= d\left(\frac{M'}{ML} \omega^{\hat{z}} \right) = d\left(\frac{M'}{L} dz \right) = \frac{\dot{M}'L - M'\dot{L}}{L^2} dt \wedge dz + \frac{M''L - M'L'}{L^2} dr \wedge dz \\
 &= \frac{M'\dot{L} - \dot{M}'L}{ML^2} \omega^{\hat{z}} \wedge \omega^{\hat{t}} + \frac{M'L' - M''L}{ML^3} \omega^{\hat{z}} \wedge \omega^{\hat{r}} \\
 \Gamma^{\hat{z}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{r}} &= \Gamma^{\hat{z}}_{\hat{r}} \wedge \Gamma^{\hat{c}}_{\hat{r}} + \Gamma^{\hat{z}}_{\hat{t}} \wedge \Gamma^{\hat{c}}_{\hat{r}} + \Gamma^{\hat{z}}_{\hat{\phi}} \wedge \Gamma^{\hat{c}}_{\hat{r}} + \Gamma^{\hat{z}}_{\hat{z}} \wedge \Gamma^{\hat{c}}_{\hat{r}} = \Gamma^{\hat{z}}_{\hat{t}} \wedge \Gamma^{\hat{c}}_{\hat{r}} = \frac{\dot{M}\dot{L}}{ML} \omega^{\hat{z}} \wedge \omega^{\hat{r}} \\
 \Rightarrow \Omega^{\hat{z}}_{\hat{r}} &= \frac{M'\dot{L} - \dot{M}'L}{ML^2} \omega^{\hat{z}} \wedge \omega^{\hat{t}} + \frac{M'L' - M''L}{ML^3} \omega^{\hat{z}} \wedge \omega^{\hat{r}} + \frac{\dot{M}\dot{L}}{ML} \omega^{\hat{z}} \wedge \omega^{\hat{r}} \\
 &= \frac{M'\dot{L} - \dot{M}'L}{ML^2} \omega^{\hat{z}} \wedge \omega^{\hat{t}} + \left(\frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} \right) \omega^{\hat{z}} \wedge \omega^{\hat{r}} \\
 d\Gamma^{\hat{z}}_{\hat{\phi}} &= 0 \\
 \Rightarrow \Omega^{\hat{z}}_{\hat{\phi}} &= \Gamma^{\hat{z}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{\phi}} = \Gamma^{\hat{z}}_{\hat{t}} \wedge \Gamma^{\hat{c}}_{\hat{\phi}} + \Gamma^{\hat{z}}_{\hat{r}} \wedge \Gamma^{\hat{c}}_{\hat{\phi}} + \Gamma^{\hat{z}}_{\hat{\phi}} \wedge \Gamma^{\hat{c}}_{\hat{\phi}} + \Gamma^{\hat{z}}_{\hat{z}} \wedge \Gamma^{\hat{c}}_{\hat{\phi}} \\
 &= \Gamma^{\hat{z}}_{\hat{t}} \wedge \Gamma^{\hat{c}}_{\hat{\phi}} + \Gamma^{\hat{z}}_{\hat{r}} \wedge \Gamma^{\hat{c}}_{\hat{\phi}} = \frac{\dot{M}}{M} \omega^{\hat{z}} \wedge \frac{\dot{B}}{B} \omega^{\hat{\phi}} + \frac{M'}{LM} \omega^{\hat{z}} \wedge \left(-\frac{B'}{LB} \omega^{\hat{\phi}} \right) \\
 &= \left(\frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \right) \omega^{\hat{z}} \wedge \omega^{\hat{\phi}}
 \end{aligned}$$

Summarized in a matrix:

$$\Omega^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 0 & -\frac{\ddot{L}}{L} \omega^{\hat{r}} \wedge \omega^{\hat{t}} & \begin{bmatrix} -\frac{\ddot{B}}{B} \omega^{\hat{\phi}} \wedge \omega^{\hat{t}} \\ + \left(\frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL} \right) \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \end{bmatrix} & \begin{bmatrix} -\frac{\ddot{M}}{M} \omega^{\hat{z}} \wedge \omega^{\hat{t}} \\ + \left(\frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} \right) \omega^{\hat{z}} \wedge \omega^{\hat{r}} \end{bmatrix} \\ S & 0 & \begin{bmatrix} \left(\frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL} \right) \omega^{\hat{\phi}} \wedge \omega^{\hat{t}} \\ + \left(\frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} \right) \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \end{bmatrix} & \begin{bmatrix} \left(\frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} \right) \omega^{\hat{z}} \wedge \omega^{\hat{t}} \\ + \left(\frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} \right) \omega^{\hat{z}} \wedge \omega^{\hat{r}} \end{bmatrix} \\ S & AS & 0 & \left(\frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \right) \omega^{\hat{z}} \wedge \omega^{\hat{\phi}} \\ S & AS & AS & 0 \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row.

Now we can find the independent elements of the Riemann tensor in the non-coordinate basis:

$$\begin{aligned}
 R^{\hat{r}\hat{r}\hat{t}}(A) &= -\frac{\dot{L}}{L} & R^{\hat{\phi}\hat{t}\hat{\phi}}(B) &= -\frac{\dot{B}}{B} & R^{\hat{z}\hat{z}\hat{t}}(C) &= -\frac{\dot{M}}{M} \\
 R^{\hat{\phi}\hat{t}\hat{r}}(D) &= \frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL} & R^{\hat{z}\hat{z}\hat{r}}(E) &= \frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} \\
 R^{\hat{\phi}\hat{r}\hat{r}}(F) &= \frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} & R^{\hat{z}\hat{r}\hat{r}}(G) &= \frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} \\
 & & R^{\hat{\phi}\hat{z}\hat{\phi}}(H) &= \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM}
 \end{aligned}$$

Where A,B,C,D,E,F,G,H will be used later, to make the calculations easier

7.8.4 The Ricci tensor

$$\begin{aligned}
 R_{\hat{a}\hat{b}} &= R^{\hat{c}\hat{a}\hat{c}\hat{b}} \\
 R_{\hat{t}\hat{t}} &= R^{\hat{c}\hat{t}\hat{c}\hat{t}} = R^{\hat{t}\hat{t}\hat{t}\hat{t}} + R^{\hat{r}\hat{t}\hat{r}\hat{t}} + R^{\hat{\phi}\hat{t}\hat{\phi}\hat{t}} + R^{\hat{z}\hat{t}\hat{z}\hat{t}} = R^{\hat{r}\hat{t}\hat{r}\hat{t}} + R^{\hat{\phi}\hat{t}\hat{\phi}\hat{t}} + R^{\hat{z}\hat{t}\hat{z}\hat{t}} \\
 &= -\frac{\dot{L}}{L} - \frac{\dot{B}}{B} - \frac{\dot{M}}{M} = A + B + C \\
 R_{\hat{r}\hat{r}} &= R^{\hat{c}\hat{r}\hat{c}\hat{r}} = R^{\hat{t}\hat{r}\hat{t}\hat{r}} + R^{\hat{r}\hat{r}\hat{r}\hat{r}} + R^{\hat{\phi}\hat{r}\hat{\phi}\hat{r}} + R^{\hat{z}\hat{r}\hat{z}\hat{r}} = R^{\hat{\phi}\hat{r}\hat{\phi}\hat{r}} + R^{\hat{z}\hat{r}\hat{z}\hat{r}} \\
 &= \frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL} + \frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} = D + E \\
 R_{\hat{\phi}\hat{t}} &= R^{\hat{c}\hat{\phi}\hat{c}\hat{t}} = R^{\hat{t}\hat{\phi}\hat{t}\hat{t}} + R^{\hat{r}\hat{\phi}\hat{r}\hat{t}} + R^{\hat{\phi}\hat{\phi}\hat{\phi}\hat{t}} + R^{\hat{z}\hat{\phi}\hat{z}\hat{t}} = 0 \\
 R_{\hat{z}\hat{t}} &= R^{\hat{c}\hat{z}\hat{c}\hat{t}} = R^{\hat{t}\hat{z}\hat{t}\hat{t}} + R^{\hat{r}\hat{z}\hat{r}\hat{t}} + R^{\hat{\phi}\hat{z}\hat{\phi}\hat{t}} + R^{\hat{z}\hat{z}\hat{z}\hat{t}} = 0 \\
 R_{\hat{r}\hat{r}} &= R^{\hat{c}\hat{r}\hat{c}\hat{r}} = R^{\hat{t}\hat{r}\hat{t}\hat{r}} + R^{\hat{r}\hat{r}\hat{r}\hat{r}} + R^{\hat{\phi}\hat{r}\hat{\phi}\hat{r}} + R^{\hat{z}\hat{r}\hat{z}\hat{r}} = -R^{\hat{t}\hat{r}\hat{t}\hat{r}} + R^{\hat{\phi}\hat{r}\hat{\phi}\hat{r}} + R^{\hat{z}\hat{r}\hat{z}\hat{r}} \\
 &= \frac{\dot{L}}{L} + \frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} + \frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} = -A + F + G \\
 R_{\hat{\phi}\hat{\phi}} &= R^{\hat{c}\hat{\phi}\hat{c}\hat{\phi}} = R^{\hat{t}\hat{\phi}\hat{t}\hat{\phi}} + R^{\hat{r}\hat{\phi}\hat{r}\hat{\phi}} + R^{\hat{\phi}\hat{\phi}\hat{\phi}\hat{\phi}} + R^{\hat{z}\hat{\phi}\hat{z}\hat{\phi}} = -R^{\hat{t}\hat{\phi}\hat{t}\hat{\phi}} + R^{\hat{r}\hat{\phi}\hat{r}\hat{\phi}} + R^{\hat{z}\hat{\phi}\hat{z}\hat{\phi}} \\
 &= \frac{\dot{B}}{B} + \frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} = -B + F + H \\
 R_{\hat{z}\hat{z}} &= R^{\hat{c}\hat{z}\hat{c}\hat{z}} = R^{\hat{t}\hat{z}\hat{t}\hat{z}} + R^{\hat{r}\hat{z}\hat{r}\hat{z}} + R^{\hat{\phi}\hat{z}\hat{\phi}\hat{z}} + R^{\hat{z}\hat{z}\hat{z}\hat{z}} = -R^{\hat{t}\hat{z}\hat{t}\hat{z}} + R^{\hat{r}\hat{z}\hat{r}\hat{z}} + R^{\hat{\phi}\hat{z}\hat{\phi}\hat{z}} \\
 &= \frac{\dot{M}}{M} + \frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} = -C + G + H
 \end{aligned}$$

Summarized in a matrix:

$$R_{\hat{a}\hat{b}} = \begin{pmatrix} -\frac{\ddot{L}}{L} - \frac{\ddot{B}}{B} - \frac{\ddot{M}}{M} & \frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL} + \frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} & 0 & 0 \\ S & \left[\frac{\ddot{L}}{L} + \frac{B'L' - B''L}{BL^3} + \frac{\dot{M}\dot{L}}{ML} \right] & 0 & 0 \\ 0 & 0 & \left[\frac{\ddot{B}}{B} + \frac{B'L' - B''L}{BL^3} \right] & 0 \\ 0 & 0 & \left[+\frac{\dot{B}\dot{L}}{BL} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \right] & 0 \\ 0 & 0 & 0 & \left[\frac{\ddot{M}}{M} + \frac{M'L' - M''L}{ML^3} \right] \\ & & & \left[+\frac{\dot{M}\dot{L}}{ML} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \right] \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row

7.8.5 The Ricci scalar

$$\begin{aligned} R &= \eta^{\hat{a}\hat{b}} R_{\hat{a}\hat{b}} \\ R &= \eta^{\hat{t}\hat{t}} R_{\hat{t}\hat{t}} + \eta^{\hat{r}\hat{r}} R_{\hat{r}\hat{r}} + \eta^{\hat{\phi}\hat{\phi}} R_{\hat{\phi}\hat{\phi}} + \eta^{\hat{z}\hat{z}} R_{\hat{z}\hat{z}} = -R_{\hat{t}\hat{t}} + R_{\hat{r}\hat{r}} + R_{\hat{\phi}\hat{\phi}} + R_{\hat{z}\hat{z}} \\ &= -(A + B + C) - A + F + G - B + F + H - C + G + H \\ &= -2A - 2B - 2C + 2F + 2G + 2H \\ &= -2R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} - 2R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}} - 2R^{\hat{z}}_{\hat{t}\hat{z}\hat{t}} + 2R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}} + 2R^{\hat{z}}_{\hat{r}\hat{z}\hat{r}} + 2R^{\hat{z}}_{\hat{\phi}\hat{z}\hat{\phi}} \\ &= 2 \left(\frac{\ddot{L}}{L} + \frac{\ddot{B}}{B} + \frac{\ddot{M}}{M} + \frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} + \frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \right) \end{aligned}$$

7.8.6 The Einstein tensor

$$\begin{aligned} G_{\hat{a}\hat{b}} &= R_{\hat{a}\hat{b}} - \frac{1}{2} \eta_{\hat{a}\hat{b}} R \\ G_{\hat{t}\hat{t}} &= R_{\hat{t}\hat{t}} - \frac{1}{2} \eta_{\hat{t}\hat{t}} R = R_{\hat{t}\hat{t}} + \frac{1}{2} R = A + B + C + \frac{1}{2} (-2A - 2B - 2C + 2F + 2G + 2H) \\ &= F + G + H = R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}} + R^{\hat{z}}_{\hat{r}\hat{z}\hat{r}} + R^{\hat{z}}_{\hat{\phi}\hat{z}\hat{\phi}} \\ &= \frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} + \frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \\ G_{\hat{r}\hat{t}} &= R_{\hat{r}\hat{t}} - \frac{1}{2} \eta_{\hat{r}\hat{t}} R = R_{\hat{r}\hat{t}} = \frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL} + \frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} \\ G_{\hat{\phi}\hat{t}} &= R_{\hat{\phi}\hat{t}} - \frac{1}{2} \eta_{\hat{\phi}\hat{t}} R = 0 \\ G_{\hat{z}\hat{t}} &= R_{\hat{z}\hat{t}} - \frac{1}{2} \eta_{\hat{z}\hat{t}} R = 0 \\ G_{\hat{r}\hat{r}} &= R_{\hat{r}\hat{r}} - \frac{1}{2} \eta_{\hat{r}\hat{r}} R = R_{\hat{r}\hat{r}} - \frac{1}{2} R \\ &= -A + F + G - \frac{1}{2} (-2A - 2B - 2C + 2F + 2G + 2H) = B + C + H \\ &= R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}} + R^{\hat{z}}_{\hat{t}\hat{z}\hat{t}} + R^{\hat{z}}_{\hat{\phi}\hat{z}\hat{\phi}} = -\frac{\ddot{B}}{B} - \frac{\ddot{M}}{M} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \\ G_{\hat{\phi}\hat{r}} &= R_{\hat{\phi}\hat{r}} - \frac{1}{2} \eta_{\hat{\phi}\hat{r}} R = 0 \end{aligned}$$

$$\begin{aligned}
 G_{z\hat{r}} &= R_{z\hat{r}} - \frac{1}{2}\eta_{z\hat{r}}R = 0 \\
 G_{\hat{\phi}\hat{\phi}} &= R_{\hat{\phi}\hat{\phi}} - \frac{1}{2}\eta_{\hat{\phi}\hat{\phi}}R = R_{\hat{\phi}\hat{\phi}} - \frac{1}{2}R \\
 &= -B + F + H - \frac{1}{2}(-2A - 2B - 2C + 2F + 2G + 2H) \\
 &= A + C + G = R^{\hat{r}}_{\hat{t}\hat{t}} + R^{\hat{z}}_{\hat{t}\hat{t}} + R^{\hat{z}}_{\hat{r}\hat{r}} = -\frac{\ddot{L}}{L} - \frac{\ddot{M}}{M} + \frac{M'L' - M''L}{ML^3} + \frac{\dot{M}\dot{L}}{ML} \\
 G_{z\hat{\phi}} &= R_{z\hat{\phi}} - \frac{1}{2}\eta_{z\hat{\phi}}R = 0 \\
 G_{z\hat{z}} &= R_{z\hat{z}} - \frac{1}{2}\eta_{z\hat{z}}R = R_{z\hat{z}} - \frac{1}{2}R = -C + G + H - \frac{1}{2}(-2A - 2B - 2C + 2F + 2G + 2H) \\
 &= A + B + F = R^{\hat{r}}_{\hat{t}\hat{t}} + R^{\hat{\phi}}_{\hat{t}\hat{t}} + R^{\hat{\phi}}_{\hat{r}\hat{r}} = -\frac{\ddot{L}}{L} - \frac{\ddot{B}}{B} + \frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL}
 \end{aligned}$$

Summarized in a matrix:

$$G_{\hat{a}\hat{b}} = \left(\begin{array}{cccc}
 \left[\begin{array}{ccc} \frac{B'L' - B''L}{BL^3} + \frac{\dot{B}\dot{L}}{BL} + \frac{M'L' - M''L}{ML^3} \\ + \frac{\dot{M}\dot{L}}{ML} + \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \end{array} \right] & \left[\begin{array}{cc} \frac{B'\dot{L}}{BL^2} - \frac{\dot{B}'}{BL} \\ + \frac{M'\dot{L}}{ML^2} - \frac{\dot{M}'}{ML} \end{array} \right] & 0 & 0 \\
 S & \left[\begin{array}{cc} -\frac{\ddot{B}}{B} - \frac{\ddot{M}}{M} \\ \frac{\dot{B}\dot{M}}{BM} - \frac{B'M'}{L^2BM} \end{array} \right] & 0 & 0 \\
 0 & 0 & \left[\begin{array}{c} -\frac{\ddot{L}}{L} - \frac{\ddot{M}}{M} + \frac{\dot{M}\dot{L}}{ML} \\ + \frac{M'L' - M''L}{ML^3} \end{array} \right] & 0 \\
 0 & 0 & 0 & \left[\begin{array}{c} -\frac{\ddot{L}}{L} - \frac{\ddot{B}}{B} + \frac{\dot{B}\dot{L}}{BL} \\ + \frac{B'L' - B''L}{BL^3} \end{array} \right]
 \end{array} \right)$$

Where \hat{a} refers to column and \hat{b} to row

7.9 The Curvature of the Poincaré Half-Plane and the Poincaré Metric.

These examples emerged from mathematicians work in the 17th century and the discovery of the hyperbolic geometry and space with negative curvature.

7.9.1 The Poincaré Half-Plane

The line element

$$ds^2 = \frac{1}{y^2}dx^2 + \frac{1}{y^2}dy^2$$

The Basis One-forms

$$\omega^x = \frac{1}{y}dx \quad dx = y\omega^x$$

$$\omega^y = \frac{1}{y}dy \quad dy = y\omega^y$$

$$\eta^{ij} = \left\{ \begin{array}{cc} 1 & \\ & 1 \end{array} \right\}$$

The Ricci rotation Coefficients

$$d\omega^{\hat{a}} = -\Gamma^{\hat{a}}_{\hat{b}} \wedge \omega^{\hat{b}}$$

$$d\omega^{\hat{x}} = d\left(\frac{1}{y} dx\right) = -\frac{1}{y^2} dy \wedge dx = -\omega^{\hat{y}} \wedge \omega^{\hat{x}} = \omega^{\hat{x}} \wedge \omega^{\hat{y}}$$

$$d\omega^{\hat{y}} = 0$$

$$\Rightarrow \Gamma^{\hat{x}}_{\hat{y}} = -\omega^{\hat{x}}$$

The Curvature Two-forms and The Riemann Tensor

$$\Omega^{\hat{a}}_{\hat{b}} = d\Gamma^{\hat{a}}_{\hat{b}} + \Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} = \frac{1}{2} R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}} \omega^{\hat{c}} \wedge \omega^{\hat{d}}$$

$$d\Gamma^{\hat{x}}_{\hat{y}} = d(-\omega^{\hat{x}}) = -\omega^{\hat{x}} \wedge \omega^{\hat{y}}$$

$$\Rightarrow \Omega^{\hat{x}}_{\hat{y}} = -\omega^{\hat{x}} \wedge \omega^{\hat{y}}$$

$$\Rightarrow R^{\hat{x}}_{\hat{y}\hat{x}\hat{y}} = -1$$

The Ricci Tensor

$$R_{\hat{a}\hat{b}} = R^{\hat{c}}_{\hat{a}\hat{c}\hat{b}}$$

$$R_{\hat{x}\hat{x}} = R^{\hat{y}}_{\hat{x}\hat{y}\hat{x}} = -1$$

$$R_{\hat{y}\hat{y}} = R^{\hat{x}}_{\hat{y}\hat{x}\hat{y}} = -1$$

The Ricci scalar

$$R = R^{\hat{a}}_{\hat{a}} = R^{\hat{x}}_{\hat{x}} + R^{\hat{y}}_{\hat{y}} = -2$$

7.9.2 The Poincaré metric

The line element

$$ds^2 = \frac{4}{1-x^2-y^2} dx^2 + \frac{4}{1-x^2-y^2} dy^2 \tag{7.1.}$$

The Basis One-forms

$$\omega^{\hat{x}} = \frac{2}{\sqrt{1-x^2-y^2}} dx \quad dx = \frac{1}{2} \sqrt{1-x^2-y^2} \omega^{\hat{x}}$$

$$\omega^{\hat{y}} = \frac{2}{\sqrt{1-x^2-y^2}} dy \quad dy = \frac{1}{2} \sqrt{1-x^2-y^2} \omega^{\hat{y}}$$

$$\eta^{ij} = \begin{Bmatrix} 1 & \\ & 1 \end{Bmatrix}$$

The Ricci rotation Coefficients

$$d\omega^{\hat{a}} = -\Gamma^{\hat{a}}_{\hat{b}} \wedge \omega^{\hat{b}}$$

$$d\omega^{\hat{x}} = d\left(\frac{2}{\sqrt{1-x^2-y^2}} dx\right) = \frac{2y}{(1-x^2-y^2)^{\frac{3}{2}}} dy \wedge dx$$

$$= \frac{2y}{(1-x^2-y^2)^{\frac{3}{2}}} \left(\frac{1}{2} \sqrt{1-x^2-y^2} \omega^{\hat{y}}\right) \wedge \left(\frac{1}{2} \sqrt{1-x^2-y^2} \omega^{\hat{x}}\right)$$

$$= \frac{y}{2\sqrt{1-x^2-y^2}} \omega^{\hat{y}} \wedge \omega^{\hat{x}} = -\frac{y}{2\sqrt{1-x^2-y^2}} \omega^{\hat{x}} \wedge \omega^{\hat{y}}$$

$$d\omega^{\hat{y}} = d\left(\frac{2}{\sqrt{1-x^2-y^2}} dy\right) = \frac{2x}{(1-x^2-y^2)^{\frac{3}{2}}} dx \wedge dy = -\frac{x}{2\sqrt{1-x^2-y^2}} \omega^{\hat{y}} \wedge \omega^{\hat{x}}$$

$$\Rightarrow \Gamma^{\hat{x}}_{\hat{y}} = \frac{y}{2\sqrt{1-x^2-y^2}} \omega^{\hat{x}} + {}^{11}\text{Antisymmetric}$$

¹¹ The curvature one-forms are antisymmetric: $\Gamma_{\hat{a}\hat{b}} = -\Gamma_{\hat{b}\hat{a}}$

$$\Gamma^{\hat{y}}_{\hat{x}} = \frac{x}{2\sqrt{1-x^2-y^2}}\omega^{\hat{y}} + {}^{12}\text{Antisymmetric}$$

$$\Rightarrow \Gamma^{\hat{x}}_{\hat{y}} = \frac{y}{2\sqrt{1-x^2-y^2}}\omega^{\hat{x}} - \frac{x}{2\sqrt{1-x^2-y^2}}\omega^{\hat{y}}$$

$$\Gamma^{\hat{y}}_{\hat{x}} = \frac{x}{2\sqrt{1-x^2-y^2}}\omega^{\hat{y}} - \frac{y}{2\sqrt{1-x^2-y^2}}\omega^{\hat{x}}$$

The Curvature Two-forms and The Riemann Tensor

$$\Omega^{\hat{a}}_{\hat{b}} = d\Gamma^{\hat{a}}_{\hat{b}} + \Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} = {}^{13}d\Gamma^{\hat{a}}_{\hat{b}} = \frac{1}{2}R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}}\omega^{\hat{c}} \wedge \omega^{\hat{d}}$$

$$\Omega^{\hat{x}}_{\hat{y}} = d\Gamma^{\hat{x}}_{\hat{y}} + \Gamma^{\hat{x}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{y}} = d\left(\frac{y}{2\sqrt{1-x^2-y^2}}\omega^{\hat{x}}\right) - d\left(\frac{x}{2\sqrt{1-x^2-y^2}}\omega^{\hat{y}}\right)$$

$$= d\left(\frac{y}{2\sqrt{1-x^2-y^2}}\frac{2}{\sqrt{1-x^2-y^2}}dx\right) - d\left(\frac{x}{2\sqrt{1-x^2-y^2}}\frac{2}{\sqrt{1-x^2-y^2}}dy\right)$$

$$= d\left(\frac{y}{1-x^2-y^2}dx\right) - d\left(\frac{x}{1-x^2-y^2}dy\right)$$

$$= \frac{1-x^2+y^2}{(1-x^2-y^2)^2}dy \wedge dx - d\frac{1+x^2-y^2}{(1-x^2-y^2)^2}dx \wedge dy = \frac{2}{(1-x^2-y^2)^2}dy \wedge dx$$

$$= \frac{2}{(1-x^2-y^2)^2}\left(\frac{1}{2}\sqrt{1-x^2-y^2}\omega^{\hat{y}}\right) \wedge \left(\frac{1}{2}\sqrt{1-x^2-y^2}\omega^{\hat{x}}\right)$$

$$= \frac{1}{2(1-x^2-y^2)}\omega^{\hat{y}} \wedge \omega^{\hat{x}}$$

$$\Rightarrow \Omega^{\hat{x}}_{\hat{y}} = -\frac{1}{2(1-x^2-y^2)}\omega^{\hat{x}} \wedge \omega^{\hat{y}}$$

$$\Rightarrow R^{\hat{x}}_{\hat{y}\hat{x}\hat{y}} = -\frac{1}{2(1-x^2-y^2)}$$

The Ricci Tensor

$$R_{\hat{a}\hat{b}} = R^{\hat{c}}_{\hat{a}\hat{c}\hat{b}}$$

$$R_{\hat{x}\hat{x}} = R^{\hat{y}}_{\hat{x}\hat{y}\hat{x}} = -\frac{1}{2(1-x^2-y^2)}$$

$$R_{\hat{y}\hat{y}} = R^{\hat{x}}_{\hat{y}\hat{x}\hat{y}} = -\frac{1}{2(1-x^2-y^2)}$$

The Ricci scalar

$$R = R^{\hat{a}}_{\hat{a}} = R^{\hat{x}}_{\hat{x}} + R^{\hat{y}}_{\hat{y}} = -\frac{1}{(1-x^2-y^2)}$$

7.10 The Tolman-Bondi- de Sitter metric (Spherical dust with a cosmological constant)

7.10.1 Ricci rotation coefficients in the Tolman-Bondi- de Sitter metric

The line element:

$$ds^2 = dt^2 - e^{-2\psi(t,r)}dr^2 - R^2(t,r)d\theta^2 - R^2(t,r)\sin^2\theta d\phi^2$$

The Basis one forms

$$\omega^{\hat{t}} = dt$$

$$\omega^{\hat{r}} = e^{-\psi(t,r)}dr$$

$$\omega^{\hat{\theta}} = R(t,r)d\theta$$

$$dr = e^{\psi(t,r)}\omega^{\hat{r}}$$

$$d\theta = \frac{1}{R(2t,r)}\omega^{\hat{\theta}}$$

¹² The curvature one-forms are antisymmetric: $\Gamma_{\hat{a}\hat{b}} = -\Gamma_{\hat{b}\hat{a}}$

¹³ $\Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} = 0$

$$\omega^{\hat{\phi}} = R(t, r) \sin \theta d\phi \qquad d\phi = \frac{1}{R(t, r) \sin \theta} \omega^{\hat{\phi}}$$

$$\eta^{ij} = \begin{pmatrix} 1 & & & \\ & -1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix}$$

$$\Lambda^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 1 & & & \\ e^{-\psi(t,r)} & & & \\ & R(t, r) & & \\ & & R(t, r) \sin \theta & \end{pmatrix}$$

$$\Lambda^{\hat{b}}_{\hat{a}} = \begin{pmatrix} 1 & & & \\ e^{\psi(t,r)} & & & \\ & \frac{1}{R(t, r)} & & \\ & & \frac{1}{R(t, r) \sin \theta} & \end{pmatrix}$$

Cartan's First Structure equation and the calculation of the Ricci rotation coefficients $\Gamma^{\hat{a}}_{\hat{b}\hat{c}}$:

$$\begin{aligned} d\omega^{\hat{a}} &= -\Gamma^{\hat{a}}_{\hat{b}} \wedge \omega^{\hat{b}} \\ \Gamma^{\hat{a}}_{\hat{b}} &= \Gamma^{\hat{a}}_{\hat{b}\hat{c}} \omega^{\hat{c}} \\ \Rightarrow d\omega^{\hat{t}} &= 0 \\ d\omega^{\hat{r}} &= d(e^{-\psi(t,r)} dr) = -\dot{\psi} e^{-\psi(t,r)} dt \wedge dr = -\dot{\psi} \omega^{\hat{t}} \wedge \omega^{\hat{r}} \\ d\omega^{\hat{\theta}} &= d(R(t, r) d\theta) = \dot{R} dt \wedge d\theta + R' dr \wedge d\theta = \frac{\dot{R}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\theta}} + \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{r}} \wedge \omega^{\hat{\theta}} \\ d\omega^{\hat{\phi}} &= d(R(t, r) \sin \theta d\phi) \\ &= \dot{R} \sin \theta dt \wedge d\phi + R' \sin \theta dr \wedge d\phi + R(t, r) \cos \theta d\theta \wedge d\phi \\ &= \frac{\dot{R}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{r}} \wedge \omega^{\hat{\phi}} + \frac{\cot \theta}{R} \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \end{aligned}$$

Summarizing the curvature one forms in a matrix:

$$\Gamma^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 0 & -\dot{\psi} \omega^{\hat{r}} & \frac{\dot{R}}{R} \omega^{\hat{\theta}} & \frac{\dot{R}}{R} \omega^{\hat{\phi}} \\ -\dot{\psi} \omega^{\hat{r}} & 0 & \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\theta}} & \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\phi}} \\ \frac{\dot{R}}{R} \omega^{\hat{\theta}} & -\frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\theta}} & 0 & \frac{\cot \theta}{R} \omega^{\hat{\phi}} \\ \frac{\dot{R}}{R} \omega^{\hat{\phi}} & -\frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\phi}} & -\frac{\cot \theta}{R} \omega^{\hat{\phi}} & 0 \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row.

Now we can read off the Ricci rotation coefficients

$$\begin{aligned} \Gamma^{\hat{t}}_{\hat{r}\hat{r}} &= -\dot{\psi} & \Gamma^{\hat{r}}_{\hat{t}\hat{r}} &= -\dot{\psi} \\ \Gamma^{\hat{t}}_{\hat{\theta}\hat{\theta}} &= \frac{\dot{R}}{R} & \Gamma^{\hat{r}}_{\hat{\theta}\hat{\theta}} &= -\frac{R'}{R} e^{\psi(t,r)} \\ \Gamma^{\hat{t}}_{\hat{\phi}\hat{\phi}} &= \frac{\dot{R}}{R} & \Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}} &= -\frac{R'}{R} e^{\psi(t,r)} \\ \Gamma^{\hat{\theta}}_{\hat{t}\hat{\theta}} &= \frac{\dot{R}}{R} & \Gamma^{\hat{\phi}}_{\hat{t}\hat{\phi}} &= \frac{\dot{R}}{R} \\ \Gamma^{\hat{\theta}}_{\hat{r}\hat{\theta}} &= \frac{R'}{R} e^{\psi(t,r)} & \Gamma^{\hat{\phi}}_{\hat{r}\hat{\phi}} &= \frac{R'}{R} e^{\psi(t,r)} \end{aligned}$$

$$\Gamma^{\hat{\theta}}_{\hat{\phi}\hat{\phi}} = -\frac{\cot \theta}{R} \quad \Gamma^{\hat{\phi}}_{\hat{\theta}\hat{\phi}} = \frac{\cot \theta}{R}$$

Transformation of the Ricci rotation coefficients $\Gamma^{\hat{a}}_{\hat{b}\hat{c}}$ into the Christoffel symbols Γ^a_{bc}

We have the transformation

$$\begin{aligned} \Gamma^a_{bc} &= \Lambda^a_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{b}} \Lambda^{\hat{f}}_{\hat{c}} \\ \Rightarrow \Gamma^r_{tr} &= \Lambda^r_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{t}} \Lambda^{\hat{f}}_{\hat{r}} = \Gamma^{\hat{r}}_{\hat{t}\hat{r}} \Lambda^{\hat{t}}_{\hat{t}} = -\dot{\psi} \cdot 1 = -\dot{\psi} \\ \Gamma^t_{rr} &= \Lambda^t_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{r}} \Lambda^{\hat{f}}_{\hat{r}} = \Lambda^t_{\hat{t}} \Gamma^{\hat{t}}_{\hat{r}\hat{r}} (\Lambda^{\hat{r}}_{\hat{r}})^2 = 1(-\dot{\psi})(e^{-\psi(t,r)})^2 = -\dot{\psi} e^{-2\psi(t,r)} \\ \Gamma^{\theta}_{t\theta} &= \Lambda^{\theta}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{t}} \Lambda^{\hat{f}}_{\hat{\theta}} = \Gamma^{\hat{\theta}}_{\hat{t}\hat{\theta}} \Lambda^{\hat{t}}_{\hat{t}} = \frac{\dot{R}}{R} \cdot 1 = \frac{\dot{R}}{R} \\ \Gamma^t_{\theta\theta} &= \Lambda^t_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{\theta}} \Lambda^{\hat{f}}_{\hat{\theta}} = \Lambda^t_{\hat{t}} \Gamma^{\hat{t}}_{\hat{\theta}\hat{\theta}} (\Lambda^{\hat{\theta}}_{\hat{\theta}})^2 = 1 \cdot \frac{\dot{R}}{R} R^2 = R\dot{R} \\ \Gamma^{\theta}_{r\theta} &= \Lambda^{\theta}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{r}} \Lambda^{\hat{f}}_{\hat{\theta}} = \Gamma^{\hat{\theta}}_{\hat{r}\hat{\theta}} \Lambda^{\hat{r}}_{\hat{r}} = \frac{R'}{R} e^{\psi(t,r)} e^{-\psi(t,r)} = \frac{R'}{R} \\ \Gamma^r_{\theta\theta} &= \Lambda^r_{\hat{r}} \Gamma^{\hat{r}}_{\hat{\theta}\hat{\theta}} (\Lambda^{\hat{\theta}}_{\hat{\theta}})^2 = e^{\psi(t,r)} \left(-\frac{R'}{R} e^{\psi(t,r)} \right) R^2 = -RR' e^{2\psi(t,r)} \\ \Gamma^{\phi}_{t\phi} &= \Lambda^{\phi}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{t}} \Lambda^{\hat{f}}_{\hat{\phi}} = \Gamma^{\hat{\phi}}_{\hat{t}\hat{\phi}} \Lambda^{\hat{t}}_{\hat{t}} = \frac{\dot{R}}{R} \cdot 1 = \frac{\dot{R}}{R} \\ \Gamma^t_{\phi\phi} &= \Lambda^t_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{\phi}} \Lambda^{\hat{f}}_{\hat{\phi}} = \Lambda^t_{\hat{t}} \Gamma^{\hat{t}}_{\hat{\phi}\hat{\phi}} (\Lambda^{\hat{\phi}}_{\hat{\phi}})^2 = 1 \cdot \frac{\dot{R}}{R} R^2 \sin^2 \theta = R\dot{R} \sin^2 \theta \\ \Gamma^{\phi}_{r\phi} &= \Lambda^{\phi}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{r}} \Lambda^{\hat{f}}_{\hat{\phi}} = \Gamma^{\hat{\phi}}_{\hat{r}\hat{\phi}} \Lambda^{\hat{r}}_{\hat{r}} = \frac{R'}{R} e^{\psi(t,r)} e^{-\psi(t,r)} = \frac{R'}{R} \\ \Gamma^r_{\phi\phi} &= \Lambda^r_{\hat{r}} \Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}} (\Lambda^{\hat{\phi}}_{\hat{\phi}})^2 = e^{\psi(t,r)} \left(-\frac{R'}{R} e^{\psi(t,r)} \right) R^2 \sin^2 \theta = -RR' e^{2\psi(t,r)} \sin^2 \theta \\ \Gamma^{\phi}_{\theta\phi} &= \Lambda^{\phi}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{\theta}} \Lambda^{\hat{f}}_{\hat{\phi}} = \Gamma^{\hat{\phi}}_{\hat{\theta}\hat{\phi}} \Lambda^{\hat{\theta}}_{\hat{\theta}} = \frac{\cot \theta}{R} R = \cot \theta \\ \Gamma^{\theta}_{\phi\phi} &= \Lambda^{\theta}_{\hat{d}} \Gamma^{\hat{d}}_{\hat{e}\hat{f}} \Lambda^{\hat{e}}_{\hat{\phi}} \Lambda^{\hat{f}}_{\hat{\phi}} = \frac{1}{R} \left(-\frac{\cot \theta}{R} \right) R^2 \sin^2 \theta = -\cos \theta \sin \theta \end{aligned}$$

However by this method we do not obtain $\Gamma^r_{rr} = -\psi'$, because the Ricci rotation coefficient $\Gamma^{\hat{r}}_{\hat{r}\hat{r}}$ does not exist. The conclusion must be, that this is not a reliable method to calculate Christoffel symbols.

To check we calculate the Christoffel symbols directly from the metric

The line element:

$$ds^2 = dt^2 - e^{-2\psi(t,r)} dr^2 - R^2(t,r) d\theta^2 - R^2(t,r) \sin^2 \theta d\phi^2$$

The metric tensor and its inverse

$$g_{ab} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -e^{-2\psi(t,r)} & 0 & 0 \\ 0 & 0 & -R^2(t,r) & 0 \\ 0 & 0 & 0 & -R^2(t,r) \sin^2 \theta \end{pmatrix}$$

$$g^{ab} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -e^{2\psi(t,r)} & 0 & 0 \\ 0 & 0 & -\frac{1}{R^2(t,r)} & 0 \\ 0 & 0 & 0 & -\frac{1}{R^2(t,r) \sin^2 \theta} \end{pmatrix}$$

$$\Gamma_{abc} = \frac{1}{2} (\partial_c g_{ab} + \partial_b g_{ac} - \partial_a g_{bc})$$

$$\Gamma^a_{bc} = g^{ad} \Gamma_{dbc}$$

$$\Gamma^r_{rr} = g^{rr} \Gamma_{rrr} = g^{rr} \frac{1}{2} \partial_r (-e^{-2\psi(t,r)}) = g^{rr} \psi' e^{-2\psi(t,r)} = -e^{2\psi(t,r)} \psi' e^{-2\psi(t,r)} = -\psi'$$

$$\begin{aligned}
 \Gamma^t_{rr} &= g^{tt}\Gamma_{trr} \\
 &= g^{tt}\left(-\frac{1}{2}\partial_t(-e^{-2\psi(t,r)})\right) = g^{tt}(-\dot{\psi}e^{-2\psi(t,r)}) = 1(-\dot{\psi}e^{-2\psi(t,r)}) = -\dot{\psi}e^{-2\psi(t,r)} \\
 \Gamma^r_{tr} &= \Gamma^r_{rt} = g^{rr}\Gamma_{trr} = g^{rr}\frac{1}{2}\partial_t(-e^{-2\psi(t,r)}) = g^{rr}\dot{\psi}e^{-2\psi(t,r)} = -e^{2\psi(t,r)}\dot{\psi}e^{-2\psi(t,r)} \\
 &= -\dot{\psi} \\
 \Gamma^t_{\theta\theta} &= g^{tt}\Gamma_{t\theta\theta} = g^{tt}\left(-\frac{1}{2}\partial_t(-R^2(t,r))\right) = g^{tt}\dot{R}R = 1 \cdot \dot{R}R = \dot{R}R \\
 \Gamma^\theta_{t\theta} &= \Gamma^\theta_{\theta t} = g^{\theta\theta}\Gamma_{\theta t\theta} = g^{\theta\theta}\frac{1}{2}\partial_t(-R^2(t,r)) = g^{\theta\theta}(-\dot{R}R) = -\frac{1}{R^2}(-\dot{R}R) = \frac{\dot{R}}{R} \\
 \Gamma^r_{\theta\theta} &= g^{rr}\Gamma_{r\theta\theta} = g^{rr}\left(-\frac{1}{2}\partial_r(-R^2(t,r))\right) = g^{rr}R'R = -e^{2\psi(t,r)}R'R \\
 \Gamma^\theta_{r\theta} &= \Gamma^\theta_{\theta r} = g^{\theta\theta}\Gamma_{\theta r\theta} = g^{\theta\theta}\left(\frac{1}{2}\partial_r(-R^2(t,r))\right) = g^{\theta\theta}(-R'R) = -\frac{1}{R^2}(-R'R) = \frac{R'}{R} \\
 \Gamma^t_{\phi\phi} &= g^{tt}\Gamma_{t\phi\phi} = g^{tt}\left(-\frac{1}{2}\partial_t(-R^2(t,r)\sin^2\theta)\right) = g^{tt}\dot{R}R\sin^2\theta = 1 \cdot \dot{R}R\sin^2\theta \\
 &= \dot{R}R\sin^2\theta \\
 \Gamma^\phi_{t\phi} &= \Gamma^\phi_{\phi t} = g^{\phi\phi}\Gamma_{\phi t\phi} = g^{\phi\phi}\frac{1}{2}\partial_t(-R^2(t,r)\sin^2\theta) = g^{\phi\phi}(-\dot{R}R\sin^2\theta) \\
 &= \left(-\frac{1}{R^2\sin^2\theta}\right)(-\dot{R}R\sin^2\theta) = \frac{\dot{R}}{R} \\
 \Gamma^r_{\phi\phi} &= g^{rr}\Gamma_{r\phi\phi} = g^{rr}\left(-\frac{1}{2}\partial_r(-R^2(t,r)\sin^2\theta)\right) = g^{rr}R'R\sin^2\theta = -e^{2\psi(t,r)}R'R\sin^2\theta \\
 \Gamma^\phi_{r\phi} &= \Gamma^\phi_{\phi r} = g^{\phi\phi}\Gamma_{\phi r\phi} = g^{\phi\phi}\frac{1}{2}\partial_r(-R^2(t,r)\sin^2\theta) = g^{\phi\phi}(-R'R\sin^2\theta) \\
 &= \left(-\frac{1}{R^2\sin^2\theta}\right)(-R'R\sin^2\theta) = \frac{R'}{R} \\
 \Gamma^\theta_{\phi\phi} &= g^{\theta\theta}\Gamma_{\theta\phi\phi} = g^{\theta\theta}\left(-\frac{1}{2}\partial_\theta(-R^2(t,r)\sin^2\theta)\right) = g^{\theta\theta}R^2\sin\theta\cos\theta \\
 &= -\frac{1}{R^2}R^2\sin\theta\cos\theta = -\sin\theta\cos\theta \\
 \Gamma^\phi_{\theta\phi} &= \Gamma^\phi_{\phi\theta} = g^{\phi\phi}\Gamma_{\phi\theta\phi} = g^{\phi\phi}\frac{1}{2}\partial_\theta(-R^2(t,r)\sin^2\theta) = g^{\phi\phi}(-R^2\sin\theta\cos\theta) \\
 &= -\frac{1}{R^2\sin^2\theta}(-R^2\sin\theta\cos\theta) = \cot\theta
 \end{aligned}$$

7.10.2 The curvature two forms

$$\begin{aligned}
 \Omega^{\hat{a}}_{\hat{b}} &= d\Gamma^{\hat{a}}_{\hat{b}} + \Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} = \frac{1}{2}R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}}\omega^{\hat{c}} \wedge \omega^{\hat{d}} \\
 d\Gamma^{\hat{r}}_{\hat{t}} &= d(-\dot{\psi}\omega^{\hat{r}}) = d(-\dot{\psi}e^{-\psi(t,r)}dr) = [-\ddot{\psi}e^{-\psi(t,r)} + (\dot{\psi})^2e^{-\psi(t,r)}]dt \wedge dr \\
 &= [-\ddot{\psi} + (\dot{\psi})^2]\omega^{\hat{t}} \wedge \omega^{\hat{r}} \\
 \Gamma^{\hat{r}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} &= \Gamma^{\hat{r}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} = 0 \\
 \Rightarrow \Omega^{\hat{r}}_{\hat{t}} &= [-\ddot{\psi} + (\dot{\psi})^2]\omega^{\hat{t}} \wedge \omega^{\hat{r}} \\
 d\Gamma^{\hat{\theta}}_{\hat{t}} &= d\left(\frac{\dot{R}}{R}\omega^{\hat{\theta}}\right) = d(\dot{R}(t,r)d\theta) = \ddot{R}dt \wedge d\theta + (\dot{R})'dr \wedge d\theta \\
 &= \frac{\ddot{R}}{R}\omega^{\hat{t}} \wedge \omega^{\hat{\theta}} + \frac{(\dot{R})'}{R}e^{\psi(t,r)}\omega^{\hat{r}} \wedge \omega^{\hat{\theta}}
 \end{aligned}$$

$$\begin{aligned}
 \Gamma^{\hat{\theta}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} &= \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{\theta}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{t}} + \Gamma^{\hat{\theta}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} = \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} \\
 &= \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\theta}} \wedge (-\dot{\psi} \omega^{\hat{r}}) \\
 \Rightarrow \quad \Omega^{\hat{\theta}}_{\hat{t}} &= \frac{\ddot{R}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\theta}} + \frac{(\dot{R})' + R'\dot{\psi}}{R} e^{\psi(t,r)} \omega^{\hat{r}} \wedge \omega^{\hat{\theta}} \\
 d\Gamma^{\hat{\phi}}_{\hat{t}} &= d\left(\frac{\dot{R}}{R} \omega^{\hat{\phi}}\right) = d(\dot{R}(t,r) \sin \theta d\phi) \\
 &= \ddot{R} \sin \theta dt \wedge d\phi + (\dot{R})' \sin \theta dr \wedge d\phi + \dot{R} \cos \theta d\theta \wedge d\phi \\
 &= \frac{\ddot{R}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{(\dot{R})'}{R} e^{\psi(t,r)} \omega^{\hat{r}} \wedge \omega^{\hat{\phi}} + \frac{\dot{R}}{R^2} \cot \theta \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\
 \Gamma^{\hat{\phi}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} &= \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} = \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{t}} \\
 &= \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\phi}} \wedge (-\dot{\psi} \omega^{\hat{r}}) + \frac{\cot \theta}{R} \omega^{\hat{\phi}} \wedge \frac{\dot{R}}{R} \omega^{\hat{\theta}} \\
 \Rightarrow \quad \Omega^{\hat{\phi}}_{\hat{t}} &= \frac{\ddot{R}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{(\dot{R})' + R'\dot{\psi}}{R} e^{\psi(t,r)} \omega^{\hat{r}} \wedge \omega^{\hat{\phi}} \\
 d\Gamma^{\hat{\theta}}_{\hat{r}} &= d\left(\frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\theta}}\right) = d(R' e^{\psi(t,r)} d\theta) \\
 &= [(\dot{R})' e^{\psi(t,r)} + R'\dot{\psi} e^{\psi(t,r)}] dt \wedge d\theta + [R'' e^{\psi(t,r)} + R'\psi' e^{\psi(t,r)}] dr \wedge d\theta \\
 &= [(\dot{R})' + R'\dot{\psi}] \frac{e^{\psi(t,r)}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\theta}} + [R'' + R'\psi'] \frac{e^{2\psi(t,r)}}{R} \omega^{\hat{r}} \wedge \omega^{\hat{\theta}} \\
 \Gamma^{\hat{\theta}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{r}} &= \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{r}} + \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{r}} + \Gamma^{\hat{\theta}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{r}} + \Gamma^{\hat{\theta}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{r}} = \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{r}} \\
 &= \frac{\dot{R}}{R} \omega^{\hat{\theta}} \wedge (-\dot{\psi} \omega^{\hat{r}}) \\
 \Rightarrow \quad \Omega^{\hat{\theta}}_{\hat{r}} &= [(\dot{R})' + R'\dot{\psi}] \frac{e^{\psi(t,r)}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\theta}} + \left[(R'' + R'\psi') \frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R} \right] \omega^{\hat{r}} \wedge \omega^{\hat{\theta}} \\
 d\Gamma^{\hat{\phi}}_{\hat{r}} &= d\left(\frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\phi}}\right) = d(R' e^{\psi(t,r)} \sin \theta d\phi) \\
 &= {}^{14} [(\dot{R})' + R'\dot{\psi}] \frac{e^{\psi(t,r)}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + [R'' + R'\psi'] \frac{e^{2\psi(t,r)}}{R} \omega^{\hat{r}} \wedge \omega^{\hat{\phi}} + \frac{R'}{R^2} e^{\psi(t,r)} \cot \theta \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\
 \Gamma^{\hat{\phi}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{r}} &= \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{r}} = \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{r}} \\
 &= \frac{\dot{R}}{R} \omega^{\hat{\phi}} \wedge (-\dot{\psi} \omega^{\hat{r}}) + \frac{\cot \theta}{R} \omega^{\hat{\phi}} \wedge \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\theta}} \\
 \Rightarrow \quad \Omega^{\hat{\phi}}_{\hat{r}} &= [(\dot{R})' + R'\dot{\psi}] \frac{e^{\psi(t,r)}}{R} \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \left[(R'' + R'\psi') \frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R} \right] \omega^{\hat{r}} \wedge \omega^{\hat{\phi}} \\
 d\Gamma^{\hat{\theta}}_{\hat{\theta}} &= d\left(\frac{\cot \theta}{R} \omega^{\hat{\theta}}\right) = d(\cos \theta d\phi) = -\sin \theta d\theta \wedge d\phi = -\frac{1}{R^2} \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\
 \Gamma^{\hat{\theta}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{\theta}} &= \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{\theta}} + \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{\theta}} + \Gamma^{\hat{\theta}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{\theta}} + \Gamma^{\hat{\theta}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{\theta}} = \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{\theta}} + \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{\theta}} \\
 &= \frac{\dot{R}}{R} \omega^{\hat{\theta}} \wedge \frac{\dot{R}}{R} \omega^{\hat{\theta}} + \frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\phi}} \wedge \left(-\frac{R'}{R} e^{\psi(t,r)} \omega^{\hat{\theta}} \right) \\
 \Rightarrow \quad \Omega^{\hat{\theta}}_{\hat{\theta}} &= \left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2} e^{2\psi(t,r)} \right] \omega^{\hat{\theta}} \wedge \omega^{\hat{\theta}}
 \end{aligned}$$

¹⁴ = $\left[(\dot{R})' e^{\psi(t,r)} \sin \theta + R'\dot{\psi} e^{\psi(t,r)} \sin \theta \right] dt \wedge d\phi + \left[R'' e^{\psi(t,r)} \sin \theta + R'\psi' e^{\psi(t,r)} \sin \theta \right] dr \wedge d\phi + R' e^{\psi(t,r)} \cos \theta d\theta \wedge d\phi =$

Summarized in a matrix:

$$\Omega^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 0 & [-\dot{\psi} + (\dot{\psi})^2]\omega^{\hat{t}} \wedge \omega^{\hat{r}} & \left[\begin{array}{l} \frac{\ddot{R}}{R}\omega^{\hat{t}} \wedge \omega^{\hat{\theta}} \\ + \frac{(\dot{R})' + R'\dot{\psi}}{R}e^{\psi}\omega^{\hat{r}} \wedge \omega^{\hat{\theta}} \end{array} \right] & \left[\begin{array}{l} \frac{\ddot{R}}{R}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} \\ + \frac{(\dot{R})' + R'\dot{\psi}}{R}e^{\psi}\omega^{\hat{r}} \wedge \omega^{\hat{\phi}} \end{array} \right] \\ S & 0 & \left[\begin{array}{l} [(\dot{R})' + R'\dot{\psi}]\frac{e^{\psi}}{R}\omega^{\hat{t}} \wedge \omega^{\hat{\theta}} \\ + [(R'' + R'\psi')\frac{e^{2\psi}}{R} + \frac{\dot{R}\dot{\psi}}{R}]\omega^{\hat{r}} \wedge \omega^{\hat{\theta}} \end{array} \right] & \left[\begin{array}{l} [(\dot{R})' + R'\dot{\psi}]\frac{e^{\psi}}{R}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} \\ + [(R'' + R'\psi')\frac{e^{2\psi}}{R} + \frac{\dot{R}\dot{\psi}}{R}]\omega^{\hat{r}} \wedge \omega^{\hat{\phi}} \end{array} \right] \\ S & AS & 0 & \left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi(t,r)} \right]\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\ S & AS & AS & 0 \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row.

Now we can find the independent elements of the Riemann tensor in the non-coordinate basis:

$$\begin{aligned} R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}}(A) &= [\dot{\psi} - (\dot{\psi})^2] \\ R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}}(B) &= -\frac{\ddot{R}}{R} \\ R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{r}}(C) &= -\left[(\dot{R})' + R'\dot{\psi} \right] \frac{e^{\psi(t,r)}}{R} \\ R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}}(D) &= -\left[(R'' + R'\psi') \frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R} \right] \\ R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}}(B) &= -\frac{\ddot{R}}{R} \\ R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{r}}(C) &= -\left[(\dot{R})' + R'\dot{\psi} \right] \frac{e^{\psi(t,r)}}{R} \\ R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}}(D) &= -\left[(R'' + R'\psi') \frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R} \right] \\ R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}}(E) &= \left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi(t,r)} \right] \end{aligned}$$

Where A,B,C,D,E will be used later, to make the calculations easier.

7.10.3 The Ricci Tensor

$$\begin{aligned} R_{\hat{a}\hat{b}} &= R^{\hat{c}}_{\hat{a}\hat{c}\hat{b}} \\ \Rightarrow R_{\hat{t}\hat{t}} &= R^{\hat{c}}_{\hat{t}\hat{c}\hat{t}} = R^{\hat{t}}_{\hat{t}\hat{t}\hat{t}} + R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}} = R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}} = [\dot{\psi} - (\dot{\psi})^2] - 2\frac{\ddot{R}}{R} \\ &= A + 2B \\ R_{\hat{r}\hat{t}} &= R^{\hat{c}}_{\hat{r}\hat{c}\hat{t}} = R^{\hat{t}}_{\hat{r}\hat{t}\hat{t}} + R^{\hat{r}}_{\hat{r}\hat{r}\hat{t}} + R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{t}} = R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{t}} \\ &= -\left[(\dot{R})' + R'\dot{\psi} \right] \frac{e^{\psi(t,r)}}{R} - \left[(\dot{R})' + R'\dot{\psi} \right] \frac{e^{\psi(t,r)}}{R} = -2\left[(\dot{R})' + R'\dot{\psi} \right] \frac{e^{\psi(t,r)}}{R} = 2C \\ R_{\hat{\theta}\hat{t}} &= R^{\hat{c}}_{\hat{\theta}\hat{c}\hat{t}} = R^{\hat{t}}_{\hat{\theta}\hat{t}\hat{t}} + R^{\hat{r}}_{\hat{\theta}\hat{r}\hat{t}} + R^{\hat{\theta}}_{\hat{\theta}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{t}} = 0 \\ R_{\hat{\phi}\hat{t}} &= R^{\hat{c}}_{\hat{\phi}\hat{c}\hat{t}} = R^{\hat{t}}_{\hat{\phi}\hat{t}\hat{t}} + R^{\hat{r}}_{\hat{\phi}\hat{r}\hat{t}} + R^{\hat{\theta}}_{\hat{\phi}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{\phi}\hat{\phi}\hat{t}} = 0 \\ R_{\hat{r}\hat{r}} &= R^{\hat{c}}_{\hat{r}\hat{c}\hat{r}} = R^{\hat{t}}_{\hat{r}\hat{t}\hat{r}} + R^{\hat{r}}_{\hat{r}\hat{r}\hat{r}} + R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}} = -R^{\hat{t}}_{\hat{r}\hat{t}\hat{r}} + R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}} \end{aligned}$$

$$\begin{aligned}
 &= -\left[\ddot{\psi} - (\dot{\psi})^2\right] - 2\left[(R'' + R'\psi')\frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] = -A + 2D \\
 R_{\hat{\theta}\hat{\theta}} &= R^{\hat{c}}_{\hat{\theta}\hat{c}\hat{\theta}} = R^{\hat{t}}_{\hat{\theta}\hat{t}\hat{\theta}} + R^{\hat{r}}_{\hat{\theta}\hat{r}\hat{\theta}} + R^{\hat{\theta}}_{\hat{\theta}\hat{\theta}\hat{\theta}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} = -R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} + R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} \\
 &= \frac{\ddot{R}}{R} - \left[(R'' + R'\psi')\frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] + \left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi(t,r)}\right] = -B + D + E \\
 R_{\hat{\phi}\hat{\phi}} &= R^{\hat{c}}_{\hat{\phi}\hat{c}\hat{\phi}} = R^{\hat{t}}_{\hat{\phi}\hat{t}\hat{\phi}} + R^{\hat{r}}_{\hat{\phi}\hat{r}\hat{\phi}} + R^{\hat{\theta}}_{\hat{\phi}\hat{\theta}\hat{\phi}} + R^{\hat{\phi}}_{\hat{\phi}\hat{\phi}\hat{\phi}} = -R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}} + R^{\hat{\theta}}_{\hat{\phi}\hat{\theta}\hat{\phi}} \\
 &= \frac{\ddot{R}}{R} - \left[(R'' + R'\psi')\frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] + \left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi(t,r)}\right] = -B + D + E \\
 &= R_{\hat{\theta}\hat{\theta}}
 \end{aligned}$$

Summarized in a matrix:

$$R_{\hat{a}\hat{b}} = \begin{pmatrix} \left[\begin{array}{c} \ddot{\psi} - (\dot{\psi})^2 \\ -2\frac{\dot{R}}{R} \end{array}\right] & -2\left[(\dot{R})' + R'\dot{\psi}\right]\frac{e^{\psi}}{R} & 0 & 0 \\ S & \left[\begin{array}{c} -\left[\ddot{\psi} - (\dot{\psi})^2\right] \\ -2\left[(R'' + R'\psi')\frac{e^{2\psi}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] \end{array}\right] & 0 & 0 \\ 0 & 0 & \left[\begin{array}{c} \frac{\ddot{R}}{R} - \left[(R'' + R'\psi')\frac{e^{2\psi}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] \\ + \frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi} \end{array}\right] & 0 \\ 0 & 0 & 0 & \left[\begin{array}{c} \frac{\ddot{R}}{R} - \left[(R'' + R'\psi')\frac{e^{2\psi}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] \\ + \frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi} \end{array}\right] \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row.

7.10.4 The Ricci Scalar

$$\begin{aligned}
 R &= \eta^{\hat{a}\hat{b}}R_{\hat{a}\hat{b}} \\
 R &= \eta^{\hat{t}\hat{t}}R_{\hat{t}\hat{t}} + \eta^{\hat{r}\hat{r}}R_{\hat{r}\hat{r}} + \eta^{\hat{\theta}\hat{\theta}}R_{\hat{\theta}\hat{\theta}} + \eta^{\hat{\phi}\hat{\phi}}R_{\hat{\phi}\hat{\phi}} = R_{\hat{t}\hat{t}} - R_{\hat{r}\hat{r}} - R_{\hat{\theta}\hat{\theta}} - R_{\hat{\phi}\hat{\phi}} \\
 &= A + 2B - (-A + 2D) - 2(-B + D + E) = 2A + 4B - 4D - 2E \\
 &= 2R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + 4R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} - 4R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} - 2R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} \\
 &= 2\left[\ddot{\psi} - (\dot{\psi})^2\right] - 4\frac{\ddot{R}}{R} + 4\left[(R'' + R'\psi')\frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] - 2\left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi(t,r)}\right]
 \end{aligned}$$

7.10.5 The Einstein Tensor

$$\begin{aligned}
 G_{\hat{a}\hat{b}} &= R_{\hat{a}\hat{b}} - \frac{1}{2}\eta_{\hat{a}\hat{b}}R \\
 G_{\hat{t}\hat{t}} &= R_{\hat{t}\hat{t}} - \frac{1}{2}\eta_{\hat{t}\hat{t}}R = R_{\hat{t}\hat{t}} - \frac{1}{2}R = A + 2B - \frac{1}{2}(2A + 4B - 4D - 2E) = 2D + E \\
 &= 2R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} = -2\left[(R'' + R'\psi')\frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] + \left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi(t,r)}\right] \\
 &= \frac{1}{R^2}\left[1 - 2R\dot{R}\dot{\psi} + (\dot{R})^2 - (2RR'' + 2RR'\psi' + (R')^2)e^{2\psi(t,r)}\right]
 \end{aligned}$$

$$\begin{aligned}
 G_{\hat{r}\hat{t}} &= R_{\hat{r}\hat{t}} - \frac{1}{2}\eta_{\hat{r}\hat{t}}R = R_{\hat{r}\hat{t}} = R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{t}} = -2\left[\frac{(\dot{R})'}{R} + R'\psi\right]\frac{e^{\psi(t,r)}}{R} \\
 G_{\hat{\theta}\hat{t}} &= R_{\hat{\theta}\hat{t}} - \frac{1}{2}\eta_{\hat{\theta}\hat{t}}R = 0 \\
 G_{\hat{r}\hat{r}} &= R_{\hat{r}\hat{r}} - \frac{1}{2}\eta_{\hat{r}\hat{r}}R = R_{\hat{r}\hat{r}} + \frac{1}{2}R = -A + 2D + \frac{1}{2}(2A + 4B - 4D - 2E) = 2B - E \\
 &= 2R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} - R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} = -2\frac{\ddot{R}}{R} - \left[\frac{1}{R^2} + \frac{(\dot{R})^2}{R^2} - \frac{(R')^2}{R^2}e^{2\psi(t,r)}\right] \\
 &= \frac{1}{R^2}\left[(R')^2e^{2\psi(t,r)} - 2R\ddot{R} - 1 - (\dot{R})^2\right] \\
 G_{\hat{\theta}\hat{\theta}} &= R_{\hat{\theta}\hat{\theta}} - \frac{1}{2}\eta_{\hat{\theta}\hat{\theta}}R = R_{\hat{\theta}\hat{\theta}} + \frac{1}{2}R = -B + D + E + \frac{1}{2}(2A + 4B - 4D - 2E) = A + B - D \\
 &= R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} - R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} = \left[\ddot{\psi} - (\dot{\psi})^2\right] - \frac{\ddot{R}}{R} + \left[(R'' + R'\psi')\frac{e^{2\psi(t,r)}}{R} + \frac{\dot{R}\dot{\psi}}{R}\right] \\
 &= \left[\ddot{\psi} - (\dot{\psi})^2\right] + \frac{1}{R}\left[(R'' + R'\psi')e^{2\psi(t,r)} + \dot{R}\dot{\psi} - \ddot{R}\right] \\
 G_{\hat{\phi}\hat{\phi}} &= G_{\hat{\theta}\hat{\theta}} = \left[\ddot{\psi} - (\dot{\psi})^2\right] + \frac{1}{R}\left[(R'' + R'\psi')e^{2\psi(t,r)} + \dot{R}\dot{\psi} - \ddot{R}\right]
 \end{aligned}$$

Summarized in a matrix:

$$G_{\hat{a}\hat{b}} = \begin{pmatrix} \frac{1}{R^2}\left[1 - 2R\dot{R}\dot{\psi} + (\dot{R})^2 - (2RR'' + 2RR'\psi' + (R')^2)e^{2\psi}\right] & -2\left[\frac{(\dot{R})'}{R} + R'\psi\right]\frac{e^{\psi}}{R} & 0 & 0 \\ 0 & \frac{1}{R^2}\left[-2R\ddot{R} - 1 - (\dot{R})^2\right] & 0 & 0 \\ 0 & 0 & \left[\ddot{\psi} - (\dot{\psi})^2 + \frac{1}{R}\left[(R'' + R'\psi')e^{2\psi} + \dot{R}\dot{\psi} - \ddot{R}\right]\right] & 0 \\ 0 & 0 & 0 & \left[\ddot{\psi} - (\dot{\psi})^2 + \frac{1}{R}\left[(R'' + R'\psi')e^{2\psi} + \dot{R}\dot{\psi} - \ddot{R}\right]\right] \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row.

7.10.6 The Einstein tensor in the coordinate basis:

$$\begin{aligned}
 G_{ab} &= \Lambda^{\hat{c}}_a \Lambda^{\hat{d}}_b G_{\hat{c}\hat{d}} \\
 G_{tt} &= \Lambda^{\hat{c}}_t \Lambda^{\hat{d}}_t G_{\hat{c}\hat{d}} = (\Lambda^{\hat{t}}_t)^2 G_{\hat{t}\hat{t}} = \frac{1}{R^2}\left[1 - 2R\dot{R}\dot{\psi} + (\dot{R})^2 - (2RR'' + 2RR'\psi' + (R')^2)e^{2\psi}\right] \\
 G_{rt} &= \Lambda^{\hat{c}}_r \Lambda^{\hat{d}}_t G_{\hat{c}\hat{d}} = \Lambda^{\hat{r}}_r \Lambda^{\hat{t}}_t G_{\hat{r}\hat{t}} = -2\left[\frac{(\dot{R})'}{R} + \frac{R'\dot{\psi}}{R}\right] \\
 G_{rr} &= \Lambda^{\hat{c}}_r \Lambda^{\hat{d}}_r G_{\hat{c}\hat{d}} = (\Lambda^{\hat{r}}_r)^2 G_{\hat{r}\hat{r}} = \frac{1}{R^2}\left[(R')^2 - (2R\ddot{R} + 1 + (\dot{R})^2)e^{-2\psi}\right] \\
 G_{\theta\theta} &= \Lambda^{\hat{c}}_{\theta} \Lambda^{\hat{d}}_{\theta} G_{\hat{c}\hat{d}} = (\Lambda^{\hat{\theta}}_{\theta})^2 G_{\hat{\theta}\hat{\theta}} = R^2\left(\left[\ddot{\psi} - (\dot{\psi})^2\right] + \frac{1}{R}\left[(R'' + R'\psi')e^{2\psi} + \dot{R}\dot{\psi} - \ddot{R}\right]\right) \\
 G_{\phi\phi} &= \Lambda^{\hat{c}}_{\phi} \Lambda^{\hat{d}}_{\phi} G_{\hat{c}\hat{d}} = (\Lambda^{\hat{\phi}}_{\phi})^2 G_{\hat{\phi}\hat{\phi}} = R^2 \sin^2 \theta \left(\left[\ddot{\psi} - (\dot{\psi})^2\right] + \frac{1}{R}\left[(R'' + R'\psi')e^{2\psi} + \dot{R}\dot{\psi} - \ddot{R}\right]\right)
 \end{aligned}$$

Summarized in a matrix:

$$G_{ab} = \begin{pmatrix} \frac{1}{R^2} \left[1 - 2R\dot{R}\psi + (\dot{R})^2 \right] & -2 \left[\frac{(\dot{R})'}{R} + \frac{R'\psi}{R} \right] & 0 & 0 \\ s & \frac{1}{R^2} \left[- (2RR'' + 2RR'\psi' + (R')^2) e^{2\psi} \right] & 0 & 0 \\ 0 & 0 & R^2 \left[\begin{matrix} \dot{\psi} - (\dot{\psi})^2 \\ + \frac{1}{R} [(R'' + R'\psi')e^{2\psi} + \dot{R}\psi - \dot{R}] \end{matrix} \right] & 0 \\ 0 & 0 & 0 & R^2 \sin^2 \theta \left[\begin{matrix} \dot{\psi} - (\dot{\psi})^2 \\ + \frac{1}{R} [(R'' + R'\psi')e^{2\psi} + \dot{R}\psi - \dot{R}] \end{matrix} \right] \end{pmatrix}$$

Where a refers to column and b to row

7.11 "The Taub-Nut Spacetime

¹⁵The Taub-Nut space-time is an exact solution to the Einstein equation.

The line element

$$ds^2 = -\frac{dt^2}{U^2(t)} + (2l)^2 U^2(t) (dr + \cos \theta d\phi)^2 + V^2(t) (d\theta^2 + \sin^2 \theta d\phi^2)$$

If $U^2(t) > 0$ t is a time variable. If $U^2(t) < 0$ t is a space variable.

Defining

$$\begin{aligned} U^2(t) &= -1 + \frac{2(mt + l^2)}{t^2 + l^2} = \frac{2mt - t^2 + l^2}{t^2 + l^2} = \frac{P^2(t)}{V^2(t)} \\ V^2(t) &= t^2 + l^2 \\ P^2(t) &= 2mt - t^2 + l^2 \\ \Rightarrow \dot{V}(t) &= \frac{1}{2V(t)} \frac{d}{dt} (V^2(t)) = \frac{t}{V(t)} \\ \ddot{V}(t) &= \frac{d}{dt} \left(\frac{t}{V(t)} \right) = \frac{d}{dt} \left(\frac{t}{V(t)} \right) = \frac{V(t) - t\dot{V}(t)}{V^2(t)} = \frac{1}{V(t)} - \frac{t}{V^2(t)} \left(\frac{t}{V(t)} \right) \\ &= \frac{1}{V(t)} - \frac{t^2}{V^3(t)} \\ \dot{U}(t) &= \frac{1}{2U(t)} \frac{d}{dt} (U^2(t)) = \frac{1}{2U(t)} \frac{d}{dt} \left(\frac{P^2(t)}{V^2(t)} \right) \\ &= \frac{1}{2U(t)} \left(\frac{V^2(t) \frac{d}{dt} P^2(t) - P^2(t) \frac{d}{dt} V^2(t)}{V^4(t)} \right) = \frac{1}{U(t)} \left(\frac{(m-t)V^2(t) - tP^2(t)}{V^4(t)} \right) \\ &= \frac{1}{U(t)} \left(\frac{m-t}{V^2(t)} - \frac{tU^2(t)}{V^2(t)} \right) \\ \frac{d}{dt} (U(t)\dot{U}(t)) &= \frac{d}{dt} \left(\frac{m-t}{V^2(t)} - \frac{tU^2(t)}{V^2(t)} \right) = \frac{d}{dt} \left(\frac{m-t}{V^2(t)} \right) - \frac{d}{dt} \left(\frac{tU^2(t)}{V^2(t)} \right) \\ &= \frac{-V^2(t) - 2(m-t)V(t)\dot{V}(t)}{V^4(t)} - \frac{\frac{d}{dt} (tU^2(t))V^2(t) - 2tU^2(t)V(t)\dot{V}(t)}{V^4(t)} \\ &= -\frac{1}{V^2(t)} - \frac{2t(m-t)}{V^4(t)} - \frac{(U^2(t) + 2tU(t)\dot{U}(t))V^2(t) - 2t^2U^2(t)}{V^4(t)} \\ &= -\frac{1}{V^2(t)} - \frac{2t(m-t)}{V^4(t)} - \frac{U^2(t)}{V^2(t)} - \frac{2tV^2(t)}{V^4(t)} \left(\frac{m-t}{V^2(t)} - \frac{tU^2(t)}{V^2(t)} \right) + \frac{2t^2U^2(t)}{V^4(t)} \\ &= -\frac{1}{V^2(t)} - \frac{4t(m-t)}{V^4(t)} - \frac{U^2(t)}{V^2(t)} + \frac{4t^2U^2(t)}{V^4(t)} \end{aligned}$$

¹⁵ https://en.wikipedia.org/wiki/Taub%E2%80%93NUT_space

7.11.1 The Riemann tensor

¹⁶The TAUB-Nut space-time is a solution of the vacuum Einstein equation if the Ricci scalar $R = 0$. We choose to work in the non-coordinate basis.

The basis one forms

$$\begin{aligned}\omega^{\hat{t}} &= \frac{1}{U(t)} dt & dt &= U(t)\omega^{\hat{t}} \\ \omega^{\hat{r}} &= 2lU(t)(dr + \cos\theta d\phi) & dr &= \frac{1}{2lU(t)}\omega^{\hat{r}} - \frac{\cot\theta}{V(t)}\omega^{\hat{\phi}} \\ \omega^{\hat{\theta}} &= V(t)d\theta & d\theta &= \frac{1}{V(t)}\omega^{\hat{\theta}} \\ \omega^{\hat{\phi}} &= \sin\theta V(t)d\phi & d\phi &= \frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}} \\ \eta^{ij} &= \begin{pmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{pmatrix}\end{aligned}$$

7.11.1.1 Cartan's First Structure equation

$$\begin{aligned}d\omega^{\hat{a}} &= -\Gamma^{\hat{a}}_{\hat{b}} \wedge \omega^{\hat{b}} \\ d\omega^{\hat{t}} &= d\left(\frac{1}{U(t)} dt\right) = 0 \\ d\omega^{\hat{r}} &= -\dot{U}(t)\omega^{\hat{r}} \wedge \omega^{\hat{t}} - \frac{lU(t)}{V^2(t)}\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} + \frac{lU(t)}{V^2(t)}\omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} \\ &= d(2lU(t)dr) + d(2lU(t)\cos\theta d\phi) \\ &= {}^{17}{}_{18}{}_{19}{}_{20} - \Gamma^{\hat{r}}_{\hat{t}} \wedge \omega^{\hat{t}} - \Gamma^{\hat{r}}_{\hat{\phi}} \wedge \omega^{\hat{\phi}} - \Gamma^{\hat{r}}_{\hat{\theta}} \wedge \omega^{\hat{\theta}} \\ d\omega^{\hat{\theta}} &= d(V(t)d\theta) = \dot{V}(t)dt \wedge d\theta = -\frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\theta}} \wedge \omega^{\hat{t}} = -\Gamma^{\hat{\theta}}_{\hat{t}} \wedge \omega^{\hat{t}} \\ d\omega^{\hat{\phi}} &= d(\sin\theta V(t)d\phi) = \sin\theta \dot{V}(t)dt \wedge d\phi + \cos\theta V(t)d\theta \wedge d\phi \\ &= \frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{\cot\theta}{V(t)}\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} = -\frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\phi}} \wedge \omega^{\hat{t}} - \frac{\cot\theta}{V(t)}\omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} \\ &= -\Gamma^{\hat{\phi}}_{\hat{t}} \wedge \omega^{\hat{t}} - \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \omega^{\hat{\theta}}\end{aligned}$$

Summarizing the curvature one forms in a matrix:

¹⁶ $G_{\hat{a}\hat{b}} = R_{\hat{a}\hat{b}} - \frac{1}{2}\eta_{\hat{a}\hat{b}}R = 0 \Rightarrow \eta^{\hat{a}\hat{b}}R_{\hat{a}\hat{b}} - \frac{1}{2}\eta^{\hat{a}\hat{b}}\eta_{\hat{a}\hat{b}}R = R - \frac{1}{2}4R = 0 \Rightarrow R = 0$

¹⁷ $= 2l\dot{U}(t)dt \wedge dr + 2l\dot{U}(t)\cos\theta dt \wedge d\phi - 2lU(t)\sin\theta d\theta \wedge d\phi =$

¹⁸ $= 2l\dot{U}(t)(U(t)\omega^{\hat{t}}) \wedge \left(\frac{1}{2lU(t)}\omega^{\hat{r}} - \frac{\cot\theta}{V(t)}\omega^{\hat{\phi}}\right) + 2l\dot{U}(t)\cos\theta (U(t)\omega^{\hat{t}}) \wedge \left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) - 2lU(t)\sin\theta \left(\frac{1}{V(t)}\omega^{\hat{\theta}}\right) \wedge \left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) =$

¹⁹ $= \left(\dot{U}(t)\omega^{\hat{t}} \wedge \omega^{\hat{r}} - 2l\dot{U}(t)(U(t)\omega^{\hat{t}}) \wedge \frac{\cot\theta}{V(t)}\omega^{\hat{\phi}}\right) + 2l\dot{U}(t)\cos\theta (U(t)\omega^{\hat{t}}) \wedge \left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) -$

$2lU(t)\sin\theta \left(\frac{1}{V(t)}\omega^{\hat{\theta}}\right) \wedge \left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) =$

²⁰ $= \dot{U}(t)\omega^{\hat{t}} \wedge \omega^{\hat{r}} - \frac{2l\dot{U}(t)U(t)\cot\theta}{V(t)}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{2l\dot{U}(t)U(t)\cot\theta}{V(t)}\omega^{\hat{t}} \wedge \omega^{\hat{\phi}} - \frac{2lU(t)}{V^2(t)}\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} =$

$$\Gamma^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 0 & \dot{U}(t)\omega^{\hat{r}} & \frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\theta}} & \frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\phi}} \\ S & 0 & \frac{lU(t)}{V^2(t)}\omega^{\hat{\phi}} & -\frac{lU(t)}{V^2(t)}\omega^{\hat{\theta}} \\ S & AS & 0 & \frac{\cot\theta}{V(t)}\omega^{\hat{\phi}} \\ S & AS & AS & 0 \end{pmatrix}$$

Where \hat{a} refers to column and \hat{b} to row.

7.11.1.2 The curvature two forms

$$\begin{aligned} \Omega^{\hat{a}}_{\hat{b}} &= d\Gamma^{\hat{a}}_{\hat{b}} + \Gamma^{\hat{a}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{b}} \\ d\Gamma^{\hat{r}}_{\hat{t}} &= d(\dot{U}(t)\omega^{\hat{r}}) = d\left(\dot{U}(t)2lU(t)(dr + \cos\theta d\phi)\right) \\ &= {}^{21} {}^{22} {}^{23} \left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\omega^{\hat{t}} \wedge \omega^{\hat{r}} - 2\frac{l\dot{U}(t)U(t)}{V^2(t)}\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\ \Gamma^{\hat{r}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} &= \Gamma^{\hat{r}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{t}} + \Gamma^{\hat{r}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} \\ &= \left(-\frac{lU(t)}{V^2(t)}\omega^{\hat{\phi}}\right) \wedge \left(\frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\theta}}\right) + \left(\frac{lU(t)}{V^2(t)}\omega^{\hat{\theta}}\right) \wedge \left(\frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\phi}}\right) \\ &= -\frac{2lU^2(t)\dot{V}(t)}{V^3(t)}\omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} \\ \Rightarrow \Omega^{\hat{r}}_{\hat{t}} &= \left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\omega^{\hat{t}} \wedge \omega^{\hat{r}} - 2\frac{l\dot{U}(t)U(t)}{V^2(t)}\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} - \frac{2lU^2(t)\dot{V}(t)}{V^3(t)}\omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} \\ &= \left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\omega^{\hat{t}} \wedge \omega^{\hat{r}} - \frac{2lU(t)}{V^3(t)}\left(\dot{U}(t)V(t) - U(t)\dot{V}(t)\right)\omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\ d\Gamma^{\hat{\theta}}_{\hat{t}} &= d\left(\frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\theta}}\right) = d(U(t)\dot{V}(t)d\theta) = \left(\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)\right)dt \wedge d\theta \\ &= \frac{U(t)}{V(t)}\left(\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)\right)\omega^{\hat{t}} \wedge \omega^{\hat{\theta}} \\ \Gamma^{\hat{\theta}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} &= \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{\theta}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{t}} + \Gamma^{\hat{\theta}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} \\ &= \left(\frac{lU(t)}{V^2(t)}\omega^{\hat{\phi}}\right) \wedge \left(\dot{U}(t)\omega^{\hat{r}}\right) + \left(-\frac{\cot\theta}{V(t)}\omega^{\hat{\phi}}\right) \wedge \left(\frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\phi}}\right) \\ &= \frac{lU(t)\dot{U}(t)}{V^2(t)}\omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\ \Rightarrow \Omega^{\hat{\theta}}_{\hat{t}} &= \frac{U(t)}{V(t)}\left(\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)\right)\omega^{\hat{t}} \wedge \omega^{\hat{\theta}} + \frac{lU(t)\dot{U}(t)}{V^2(t)}\omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\ d\Gamma^{\hat{\phi}}_{\hat{t}} &= d\left(\frac{U(t)\dot{V}(t)}{V(t)}\omega^{\hat{\phi}}\right) = d(U(t)\dot{V}(t)\sin\theta d\phi) \\ &= \left(\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)\right)\sin\theta dt \wedge d\phi + U(t)\dot{V}(t)\cos\theta d\theta \wedge d\phi \end{aligned}$$

$${}^{21} = 2l\left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)dt \wedge dr + 2l\left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)dt \wedge \cos\theta d\phi - \dot{U}(t)2lU(t)\sin\theta d\theta \wedge d\phi =$$

$${}^{22} = 2l\left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\left(U(t)\omega^{\hat{t}}\right) \wedge \left(\frac{1}{2lU(t)}\omega^{\hat{r}} - \frac{\cot\theta}{V(t)}\omega^{\hat{\phi}}\right) + 2l\left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\left(U(t)\omega^{\hat{t}}\right) \wedge$$

$$\cos\theta\left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) - \dot{U}(t)2lU(t)\sin\theta\left(\frac{1}{V(t)}\omega^{\hat{\theta}}\right) \wedge \left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) =$$

$${}^{23} = 2l\left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\left(U(t)\omega^{\hat{t}}\right) \wedge \left(\frac{1}{2lU(t)}\omega^{\hat{r}}\right) - 2l\left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\left(U(t)\omega^{\hat{t}}\right) \wedge \left(\frac{\cot\theta}{V(t)}\omega^{\hat{\phi}}\right) +$$

$$2l\left(\ddot{U}(t)U(t) + \dot{U}^2(t)\right)\left(U(t)\omega^{\hat{t}}\right) \wedge \cos\theta\left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) - \dot{U}(t)2lU(t)\sin\theta\left(\frac{1}{V(t)}\omega^{\hat{\theta}}\right) \wedge \left(\frac{1}{\sin\theta V(t)}\omega^{\hat{\phi}}\right) =$$

$$\begin{aligned}
 &= \frac{U(t)}{V(t)} (\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)) \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{U(t)\dot{V}(t) \cot \theta}{V^2(t)} \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\
 \Gamma^{\hat{\phi}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{t}} &= \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{t}} + \Gamma^{\hat{\phi}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{t}} \\
 &= -\frac{lU(t)\dot{U}(t)}{V^2(t)} \omega^{\hat{\theta}} \wedge \omega^{\hat{r}} + \frac{U(t)\dot{V}(t) \cot \theta}{V^2(t)} \omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} \\
 \Rightarrow \Omega^{\hat{\phi}}_{\hat{t}} &= \frac{U(t)}{V(t)} (\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)) \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} - \frac{lU(t)\dot{U}(t)}{V^2(t)} \omega^{\hat{\theta}} \wedge \omega^{\hat{r}} \\
 d\Gamma^{\hat{\theta}}_{\hat{r}} &= d\left(\frac{lU(t)}{V^2(t)} \omega^{\hat{\phi}}\right) = d\left(\frac{lU(t)}{V(t)} \sin \theta d\phi\right) \\
 &= l\left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) \sin \theta dt \wedge d\phi + \frac{lU(t)}{V(t)} \cos \theta d\theta \wedge d\phi \\
 &= l\frac{U(t)}{V(t)} \left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{lU(t)}{V^3(t)} \cot \theta \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\
 \Gamma^{\hat{\theta}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{r}} &= \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{r}} + \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{r}} + \Gamma^{\hat{\theta}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{r}} + \Gamma^{\hat{\theta}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{r}} \\
 &= \frac{U(t)\dot{U}(t)\dot{V}(t)}{V(t)} \omega^{\hat{\theta}} \wedge \omega^{\hat{r}} + \frac{lU(t)}{V^3(t)} \cot \theta \omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} \\
 \Rightarrow \Omega^{\hat{\theta}}_{\hat{r}} &= l\frac{U(t)}{V(t)} \left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) \omega^{\hat{t}} \wedge \omega^{\hat{\phi}} + \frac{U(t)\dot{U}(t)\dot{V}(t)}{V(t)} \omega^{\hat{\theta}} \wedge \omega^{\hat{r}} \\
 d\Gamma^{\hat{\phi}}_{\hat{r}} &= d\left(-\frac{lU(t)}{V^2(t)} \omega^{\hat{\theta}}\right) = -l\left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) dt \wedge d\theta \\
 &= -l\frac{U(t)}{V(t)} \left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) \omega^{\hat{t}} \wedge \omega^{\hat{\theta}} \\
 \Gamma^{\hat{\phi}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{r}} &= \Gamma^{\hat{\phi}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{r}} + \Gamma^{\hat{\phi}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{r}} = \frac{U(t)\dot{U}(t)\dot{V}(t)}{V(t)} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\
 \Rightarrow \Omega^{\hat{\phi}}_{\hat{r}} &= -l\frac{U(t)}{V(t)} \left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) \omega^{\hat{t}} \wedge \omega^{\hat{\theta}} + \frac{U(t)\dot{U}(t)\dot{V}(t)}{V(t)} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}} \\
 d\Gamma^{\hat{\theta}}_{\hat{\theta}} &= d(\cos \theta d\phi) = -\sin \theta d\theta \wedge d\phi = -\frac{1}{V^2(t)} \omega^{\hat{\theta}} \wedge \omega^{\hat{\phi}} \\
 \Gamma^{\hat{\theta}}_{\hat{c}} \wedge \Gamma^{\hat{c}}_{\hat{\theta}} &= \Gamma^{\hat{\theta}}_{\hat{t}} \wedge \Gamma^{\hat{t}}_{\hat{\theta}} + \Gamma^{\hat{\theta}}_{\hat{r}} \wedge \Gamma^{\hat{r}}_{\hat{\theta}} + \Gamma^{\hat{\theta}}_{\hat{\theta}} \wedge \Gamma^{\hat{\theta}}_{\hat{\theta}} + \Gamma^{\hat{\theta}}_{\hat{\phi}} \wedge \Gamma^{\hat{\phi}}_{\hat{\theta}} \\
 &= \left(\frac{U(t)\dot{V}(t)}{V(t)} \omega^{\hat{\phi}}\right) \wedge \left(\frac{U(t)\dot{V}(t)}{V(t)} \omega^{\hat{\theta}}\right) + \left(-\frac{lU(t)}{V^2(t)} \omega^{\hat{\theta}}\right) \wedge \left(-\frac{lU(t)}{V^2(t)} \omega^{\hat{\phi}}\right) \\
 &= \left[\left(\frac{U(t)\dot{V}(t)}{V(t)}\right)^2 - \left(\frac{lU(t)}{V^2(t)}\right)^2\right] \omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}} \\
 \Rightarrow \Omega^{\hat{\theta}}_{\hat{\theta}} &= \left[\left(\frac{U(t)\dot{V}(t)}{V(t)}\right)^2 - \left(\frac{lU(t)}{V^2(t)}\right)^2 + \frac{1}{V^2(t)}\right] \omega^{\hat{\phi}} \wedge \omega^{\hat{\theta}}
 \end{aligned}$$

7.11.1.3 The Riemann tensor:

$$\begin{aligned}
 \Omega^{\hat{a}}_{\hat{b}} &= \frac{1}{2} R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}} \omega^{\hat{c}} \wedge \omega^{\hat{d}} \\
 R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} &= -\ddot{U}(t)U(t) - \dot{U}^2(t) & R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} &= \left(\frac{U(t)\dot{V}(t)}{V(t)}\right)^2 - \left(\frac{lU(t)}{V^2(t)}\right)^2 + \frac{1}{V^2(t)} \\
 R^{\hat{r}}_{\hat{t}\hat{\theta}\hat{\phi}} &= -\frac{2lU(t)}{V^3(t)} (\dot{U}(t)V(t) - U(t)\dot{V}(t))
 \end{aligned}$$

$$\begin{aligned}
 R^{\hat{\theta}}_{\hat{t}\hat{t}} &= -\frac{U(t)}{V(t)}(\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)) &= R^{\hat{\phi}}_{\hat{t}\hat{t}} &= -\frac{U(t)}{V(t)}(\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)) \\
 R^{\hat{\theta}}_{\hat{t}\hat{r}} &= \frac{lU(t)\dot{U}(t)}{V^2(t)} &R^{\hat{\phi}}_{\hat{t}\hat{r}} &= -\frac{lU(t)\dot{U}(t)}{V^2(t)} \\
 R^{\hat{\theta}}_{\hat{r}\hat{t}} &= l\frac{U(t)}{V(t)}\left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) &R^{\hat{\phi}}_{\hat{r}\hat{t}} &= -l\frac{U(t)}{V(t)}\left(\frac{\dot{U}(t)V(t) - U(t)\dot{V}(t)}{V^2(t)}\right) \\
 R^{\hat{\theta}}_{\hat{r}\hat{r}} &= U(t)\dot{U}(t)\frac{\dot{V}(t)}{V(t)} &R^{\hat{\phi}}_{\hat{r}\hat{r}} &= U(t)\dot{U}(t)\frac{\dot{V}(t)}{V(t)}
 \end{aligned}$$

7.11.2 The Ricci tensor

The diagonal elements of the Ricci tensor:

$$\begin{aligned}
 R_{\hat{a}\hat{b}} &= R^{\hat{c}}_{\hat{a}\hat{c}\hat{b}} \\
 R_{\hat{t}\hat{t}} &= R^{\hat{t}}_{\hat{t}\hat{t}\hat{t}} + R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}} = R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + 2R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} \\
 R_{\hat{r}\hat{r}} &= R^{\hat{t}}_{\hat{r}\hat{t}\hat{r}} + R^{\hat{r}}_{\hat{r}\hat{r}\hat{r}} + R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}} = -R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + 2R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} \\
 R_{\hat{\theta}\hat{\theta}} &= R^{\hat{t}}_{\hat{\theta}\hat{t}\hat{\theta}} + R^{\hat{r}}_{\hat{\theta}\hat{r}\hat{\theta}} + R^{\hat{\theta}}_{\hat{\theta}\hat{\theta}\hat{\theta}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} = -R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} + R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} = R_{\hat{\phi}\hat{\phi}}
 \end{aligned}$$

7.11.3 The Ricci Scalar

$$\begin{aligned}
 R &= \eta^{\hat{a}\hat{b}}R_{\hat{a}\hat{b}} = \eta^{\hat{t}\hat{t}}R_{\hat{t}\hat{t}} + \eta^{\hat{r}\hat{r}}R_{\hat{r}\hat{r}} + \eta^{\hat{\theta}\hat{\theta}}R_{\hat{\theta}\hat{\theta}} + \eta^{\hat{\phi}\hat{\phi}}R_{\hat{\phi}\hat{\phi}} = -R_{\hat{t}\hat{t}} + R_{\hat{r}\hat{r}} + R_{\hat{\theta}\hat{\theta}} + R_{\hat{\phi}\hat{\phi}} \\
 &= -\left(R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + 2R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}}\right) - R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + 2R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} - R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} + R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} - R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{t}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}} \\
 &= 2\left(-R^{\hat{r}}_{\hat{t}\hat{r}\hat{t}} + 2R^{\hat{\theta}}_{\hat{r}\hat{\theta}\hat{r}} - 2R^{\hat{\theta}}_{\hat{t}\hat{\theta}\hat{t}} + R^{\hat{\phi}}_{\hat{\theta}\hat{\phi}\hat{\theta}}\right) \\
 &= {}^{24} {}^{25} {}^{26} {}^{27} 2\left(\frac{U^2(t)}{V^2(t)} - \frac{U^2(t)}{V^4(t)}(t^2 + l^2)\right) \\
 &= 0
 \end{aligned}$$

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$$\begin{aligned}
 {}^{24} &= 2\left(\frac{d}{dt}(U(t)\dot{U}(t)) + 2U(t)\dot{U}(t)\frac{\dot{V}(t)}{V(t)} + 2\frac{U(t)}{V(t)}(\dot{U}(t)\dot{V}(t) + U(t)\ddot{V}(t)) + \left(\frac{U(t)\dot{V}(t)}{V(t)}\right)^2 - \left(\frac{lU(t)}{V^2(t)}\right)^2 + \frac{1}{V^2(t)}\right) = \\
 {}^{25} &= 2\left(\frac{d}{dt}(U(t)\dot{U}(t)) + 4U(t)\dot{U}(t)\frac{\dot{V}(t)}{V(t)} + 2U^2(t)\frac{\dot{V}(t)}{V(t)} + \left(\frac{U(t)\dot{V}(t)}{V(t)}\right)^2 - \left(\frac{lU(t)}{V^2(t)}\right)^2 + \frac{1}{V^2(t)}\right) = \\
 {}^{26} &= 2\left(-\frac{1}{V^2(t)} - \frac{4t(m-t)}{V^4(t)} - \frac{U^2(t)}{V^2(t)} + \frac{4t^2U^2(t)}{V^4(t)} + 4\left(\frac{m-t}{V^2(t)} - \frac{tU^2(t)}{V^2(t)}\right)\frac{t}{V^2(t)} + 2U^2(t)\left(\frac{1}{V^2(t)} - \frac{t^2}{V^4(t)}\right) + \left(t\frac{U(t)}{V^2(t)}\right)^2 - \left(\frac{lU(t)}{V^2(t)}\right)^2 + \frac{1}{V^2(t)}\right) = \\
 {}^{27} &= 2\left(-\frac{4t(m-t)}{V^4(t)} - \frac{U^2(t)}{V^2(t)} + \frac{4t^2U^2(t)}{V^4(t)} + 4\frac{t(m-t)}{V^4(t)} - 4\frac{t^2U^2(t)}{V^4(t)} + 2\frac{U^2(t)}{V^2(t)} - 2\frac{t^2U^2(t)}{V^4(t)} + t^2\frac{U^2(t)}{V^4(t)} - \frac{l^2U^2(t)}{V^4(t)}\right) =
 \end{aligned}$$

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^a (McMahon, 2006, s. 80), (Penrose, 2004, s. 231)

^b (McMahon, 2006, s. 80)

^c (McMahon, 2006, s. 113)

^d (McMahon, 2006, s. 87)

^e (McMahon, 2006, s. 89)

^f (McMahon, 2006, s. 113)

^g (McMahon, 2006, s. 91), (Kay, 1988, s. 75)

^h (McMahon, 2006, s. 91), (Kay, 1988, s. 102)

ⁱ (McMahon, 2006, s. 120)

^j (McMahon, 2006, s. 120)

^k (McMahon, 2006, s. 84), (Hartle, *An Introduction to Einstein's General Relativity*, 2003, s. 143, 165, 184), (Kay, 1988, s. 126)

^l (McMahon, 2006, s. 120)

^m (Hartle, *An Introduction to Einstein's General Relativity*, 2003, s. 160)

ⁿ (McMahon, 2006, s. 325)

^o (McMahon, 2006, s. 325)

^p (McMahon, 2006, s. 152-53)

^q (Penrose, 2004, s. 50), (Hartle, *Gravity - An introduction to Einstein's General Relativity*, 2003, p. 184)

^r (Penrose, 2004, s. 50)

^s (McMahon, 2006, s. 106)

^t (McMahon, 2006, s. 121)

^u (Choquet-Bruhat, 2015, s. 98), (Ellis, 1973, p. 170)