

In this file you find all sorts of misprints. There are quite many but it does not mean that the book is entirely bad. On the contrary, it gives you the opportunity to carry out a lot of calculations, and it is a good starting point. If the misprint is marked by an asterix (*) it means that you can find a more thorough calculation in my file 'lots of calculations' on my homepage <http://physicsusuan.mono.net/9035/General%20Relativity>. The list is not complete, it covers the chapters 1-13 and the final exam, but I will post more on my homepage as I get along. I have made this list entirely on my own, and should you find that I have made something wrong or missed something, which I am bound to, please do not hesitate to contact me: logik.susan@gmail.com

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Contents

p.ix 1.7 ~~Eddington-Finkelstein Coordinates~~ The Radial Null Geodesic

Chapter 1

p.2 1.20 operator to the right-hand side of (1.1), we obtain

p.6 1.6-7 clock 1 at time t' , and the reflected beam arrives at the location of clock 1 at time t_2 .

p.12 l.12 $ct' =$
 p.16 l.13 $\beta_3^2 = \frac{(\beta_1 + \beta_2)^2}{(1 + \beta_1\beta_2)^2}$
 l.15 $\beta_3 = \frac{\beta_1 + \beta_2}{1 + \beta_1\beta_2}$

Chapter 2

p.29 l.18 $Vf = (V^a e_a)f = V^a \partial_a$
 p.33 l.7 as shown in Fig. 2-5
 p.34 l.4 $ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$ (2.9)
 p.40 l.15 $g^{yx}x^2 + g^{yy} = 0$
 l.16 $g^{yy} = -x^2 g^{yx}$
 p.41 l.11 $W_x = x^2(0) + 1(1) = 1$
 l.12 $W_y = 1(0) - 1(1) = -1$
 l.13 $\Rightarrow W_a = (1, -1)$
 p.42 l.2 $S^a_b =$
 l.3 $= g^{ac} T_c^b =$
 l.4 $= g_{ae} R^e_{bcd}$
 p.44 l.12 $\partial_r = \frac{\partial x}{\partial r} \partial_x + \frac{\partial y}{\partial r} \partial_y + \frac{\partial z}{\partial r} \partial_z$

Chapter 3

p.49 l.5 U_1 . A point p that belongs to the manifold M (say in U_1)
 p.55 l.3 $Q^a_b = S^a_b - T^a_b$

Chapter 4

p.67 l.2 $= \left(\frac{\partial A^b}{\partial x^a} + \Gamma^b_{ca} A^c \right) e_b$
 l.5 $\nabla_b A^a = \frac{\partial A^a}{\partial x^b} + \Gamma^a_{bc} A^c$ (4.6)
 p.72 l.19 $\Gamma_{abc} = \frac{1}{2} \left(\frac{\partial g_{ab}}{\partial x^c} + \frac{\partial g_{ac}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^a} \right)$ (4.15)
 l.21 $\Gamma_{abc} = \frac{1}{2} \left(\frac{\partial g_{ab}}{\partial x^c} + \frac{\partial g_{ac}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^a} \right) + C_{abc} + C_{acb} - C_{bca}$
 p.73 l.7 $\Gamma_{abc} = \frac{1}{2} \left(\frac{\partial g_{ab}}{\partial x^c} + \frac{\partial g_{ac}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^a} \right)$
 l.8 $\Gamma_{bca} = \frac{1}{2} \left(\frac{\partial g_{bc}}{\partial x^a} + \frac{\partial g_{ba}}{\partial x^c} - \frac{\partial g_{ca}}{\partial x^b} \right)$
 l.10 $\Gamma_{abc} + \Gamma_{bca} = \frac{1}{2} \left(\frac{\partial g_{ab}}{\partial x^c} + \frac{\partial g_{ac}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^a} \right) + \frac{1}{2} \left(\frac{\partial g_{bc}}{\partial x^a} + \frac{\partial g_{ba}}{\partial x^c} - \frac{\partial g_{ca}}{\partial x^b} \right)$
 l.11 $= \frac{1}{2} \left(\frac{\partial g_{ac}}{\partial x^b} - \frac{\partial g_{ca}}{\partial x^b} + \frac{\partial g_{ab}}{\partial x^c} + \frac{\partial g_{ba}}{\partial x^c} - \frac{\partial g_{bc}}{\partial x^a} + \frac{\partial g_{bc}}{\partial x^a} \right)$
 l.12 $= \frac{1}{2} \left(2 \frac{\partial g_{ab}}{\partial x^c} \right) = \frac{\partial g_{ab}}{\partial x^c}$
 l.17 $\Gamma^a_{bc} = \frac{1}{2} g^{ad} \left(\frac{\partial g_{db}}{\partial x^c} + \frac{\partial g_{dc}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^d} \right)$ (4.16)

- p.75 l.2 the last term
 l.3 $\Gamma^{\phi}_{bc} = \frac{1}{2} g^{\phi\phi} \left(\frac{\partial g_{\phi b}}{\partial x^c} + \frac{\partial g_{\phi c}}{\partial x^b} \right)$
 l.5 $\Gamma^{\phi}_{\phi\theta} = \frac{1}{2} g^{\phi\phi} \left(\frac{\partial g_{\phi\phi}}{\partial \theta} + \frac{\partial g_{\phi\theta}}{\partial \phi} \right) =$
- p.76 l.3 $\Gamma_{abc} = \frac{1}{2} (\partial_c g_{ab} + \partial_b g_{ac} - \partial_a g_{bc})$
 l.16-17 Looking at $\Gamma_{abc} = \frac{1}{2} (\partial_c g_{ab} + \partial_b g_{ac} - \partial_a g_{bc})$
- p.77 (*) l.2 $\Gamma_{xxu} = \frac{1}{2} (\partial_u g_{xx} + \partial_x g_{xu} - \partial_x g_{xu}) = \frac{1}{2} \partial_u g_{xx} = \frac{1}{2} [2(1-u)] = 1-u$
 l.4 $\Gamma_{uux} = \frac{1}{2} (\partial_x g_{ux} + \partial_x g_{ux} - \partial_u g_{xx}) = -\frac{1}{2} \partial_u g_{xx} = -\frac{1}{2} [2(1-u)] = -1+u$
 l.8 $\Gamma_{uyy} = 1+u$
 l.9 $\Gamma_{yyu} = -1-u$
 l.12 their last two
 l.13 $\Gamma_{abc} = \Gamma_{acb}$
- p.80 l.6 $(a\alpha + b\beta) \wedge \gamma = a\alpha \wedge \gamma + b\beta \wedge \gamma$
- p.81 l.17 $L_V W = V^b \nabla_b W - W^b \nabla_b V$ or alternatively $(L_V W)^a = V^b \nabla_b W^a - W^b \nabla_b V^a$ (4.27)
 l.19 $L_V T_{ab} = V^c \nabla_c T_{ab} + T_{cb} \nabla_a V^c + T_{ac} \nabla_b V^c$ (4.28)
- p.82 l.20-21 This should be omitted: We can write (4.33) in the more convenient form:
 $\frac{d^2 x^a}{ds^2} + \Gamma^a_{bc} \frac{dx^b}{ds} \frac{dx^c}{ds} = 0$
- p.83 (*) l.18 $\frac{d^2 \phi}{ds^2} + \frac{2}{r} \frac{dr}{ds} \frac{d\phi}{ds} = 0$
- p.85 (*) l.13 $\Rightarrow \Gamma^{\tau}_{\xi\tau} = \frac{1}{\xi} = \Gamma^{\tau}_{\tau\xi}$
- p.86 l.6 $R_{abcd} = g_{ae} R^e_{bcd}$
 (*) l.8 $R_{abcd} = \partial_c \Gamma^e_{abd} - \partial_d \Gamma^e_{abc} - \Gamma^e_{ac} \Gamma^e_{bd} + \Gamma^e_{ad} \Gamma^e_{bc}$ (4.42)
 l.12 $-\Gamma^e_{ac} \Gamma^e_{bd} + \Gamma^e_{ad} \Gamma^e_{bc}$ (4.43)
- p.87 l.3 All together, in n dimensions, there are $n^2(n^2 - 1)/12$ independent ~~nonzero~~ components ...
 l.10 $R_{1212} ; R_{1313} ; R_{2323} ; R_{1213} ; R_{1232} = R_{2123} ; R_{1323} = R_{3132}$
- p.88 l.1 ... of the Riemann tensor in two **dimensions** are ...
- p.91 (*) Q2 $\Gamma_{\theta\phi\phi} = \frac{1}{2} (-\partial_{\theta} g_{\phi\phi}) = \frac{1}{2} (-\partial_{\theta} (r^2 \sin^2 \theta)) = -r^2 \cos \theta \sin \theta$
- p.92 (*) Q7 $\Gamma^v_{xx} = -1+u$

Chapter 5

- p.105 l.17 $\alpha \wedge \beta = (-1)^{pq} \beta \wedge \alpha$ (5.13)
- p.120 (*) l.4 $R^{\hat{\phi}}_{\hat{r}\hat{\phi}} = -\frac{\dot{a}^2 + k}{a^2}$
 (*) Q3 $\Gamma^r_{\phi\phi} = (\Lambda^{-1})^r_{\hat{r}} \Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}} \Lambda^{\hat{\phi}}_{\phi} \Lambda^{\hat{\phi}}_{\phi} = 1 \left(-\frac{1}{r} \right) (r \sin \theta)^2 = -r \sin^2 \theta$
 (*) Q4 $\Gamma^{\hat{u}}_{\hat{v}\hat{v}} = \Gamma^{\hat{v}}_{\hat{u}\hat{v}} = +\frac{1}{u}$
- p.121 Q5 Where are the "hats"?

(*) Q7
$$G_{\hat{t}\hat{t}} = \frac{1}{R^2} [-Re^{2\Psi}(2R'\Psi' + 2R'' + R^{-1}R'^2) - 2\dot{R}\dot{\Psi}R + 1 + \dot{R}^2]$$

Chapter 6

- p.130 l.15 Cases 2 and 3 in our
- p.133 l.5 $L_u \eta^a = u^b \nabla_b \eta^a - \eta^b \nabla_b u^a = 0$
- l.19 $\nabla_a S^c_b = \partial_a S^c_b + \Gamma^c_{ad} S^d_b - \Gamma^d_{ba} S^c_d$
- l.24 $\nabla_b \nabla_a V^c = \partial_b (\partial_a V^c + \dots$
- p.134 l.4 $\partial_a (\partial_b V^c - \Gamma^c_{eb} V^e) - \partial_b (\partial_a V^c + \dots$ (6.2)
- p.135 l.2 $= R^c_{\hat{d}\hat{a}\hat{b}} V^{\hat{d}}$ (6.3)
- p.136 l.1 However, since u^a is the tangent
- p.137 l.12 $\Gamma^a_{bc} = \frac{1}{2} g^{ad} \left(\frac{\partial g_{db}}{\partial x^c} + \frac{\partial g_{dc}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^d} \right)$ (4.16)
- p.139 l.21 $ds^2 = -dt^2 + e^{2b(t,r)} dr^2 + R^2(t,r) d\phi^2$
- p.141 (*) l.18 $= -\frac{1}{R} \frac{\partial R}{\partial t} \omega^{\hat{\phi}} \wedge \omega^{\hat{t}} - \frac{1}{R} \frac{\partial R}{\partial r} e^{-b(t,r)} \omega^{\hat{\phi}} \wedge \omega^{\hat{r}}$
- p.143 l.3 $\Rightarrow \Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}} =$
- l.7 $= \frac{1}{2} R^{\hat{a}}_{\hat{b}\hat{c}\hat{d}} \omega^{\hat{c}} \wedge \omega^{\hat{d}}$ (6.21)
- p.144 l.10 $\Omega^{\hat{r}}_{\hat{t}} =$
- p.146 l.4 we write out $\Omega^{\hat{t}}_{\hat{\phi}} =$
- l.6 $\Omega^{\hat{t}}_{\hat{\phi}} =$
- l.11 $R^{\hat{t}}_{\hat{\phi}\hat{t}\hat{\phi}} =$
- l.12 All together, the independent nonzero components
- p.147 l.2 Using the results of Example 6-2
- (*) l.15 $R^{\hat{\phi}}_{\hat{t}\hat{\phi}\hat{r}} = -R^{\hat{t}}_{\hat{\phi}\hat{r}\hat{\phi}}$
- l.18 $R_{\hat{r}\hat{r}} = R^{\hat{c}}_{\hat{r}\hat{c}\hat{r}} = R^{\hat{t}}_{\hat{r}\hat{t}\hat{r}} + R^{\hat{r}}_{\hat{r}\hat{r}\hat{r}} + R^{\hat{\phi}}_{\hat{r}\hat{\phi}\hat{r}}$ (6.29)
- p.149 (*) l.5 $G_{\hat{t}\hat{r}} = -\frac{e^{-b(t,r)}}{R} \frac{\partial^2 R}{\partial t \partial r} + \frac{e^{-b(t,r)}}{R} \frac{\partial b}{\partial t} \frac{\partial R}{\partial r}$
- l.8 $ds^2 = -dt^2 + e^{2b(t,r)} dr^2 + R^2(t,r) d\phi^2$, this is easy enough:
- (*) l.9 $\Lambda^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{b(t,r)} & 0 \\ 0 & 0 & R(t,r) \end{pmatrix}$ (6.33)
- (*) l.15 $G_{\hat{t}\hat{t}} = \Lambda^{\hat{t}}_{\hat{t}} \Lambda^{\hat{t}}_{\hat{t}} G_{\hat{t}\hat{t}} = (1)(1)G_{\hat{t}\hat{t}}$ (6.35)
- (*) l.18 $(1)(e^{b(t,r)}) \left(-\frac{e^{-b(t,r)}}{R} \frac{\partial^2 R}{\partial t \partial r} + \frac{e^{-b(t,r)}}{R} \frac{\partial b}{\partial t} \frac{\partial R}{\partial r} \right)$ (6.36)
- p.153 l.10 $\Gamma^{\hat{t}}_{\hat{r}\hat{r}}$
- l.15 $\Gamma^{\hat{r}}_{\hat{\phi}\hat{\phi}}$

Chapter 7

- p.155 l.9 tensor and energy-momentum tensor interchangeable.
- p.159 l.4 $\gamma = \frac{1}{\sqrt{1 - \left(\frac{v}{c}\right)^2}}$

- l.16 $= \gamma^2 \frac{\partial \rho}{\partial t} + \gamma^2 \frac{\partial(\rho u^x)}{\partial x} + \gamma^2 \frac{\partial(\rho u^y)}{\partial y} + \gamma^2 \frac{\partial(\rho u^z)}{\partial z}$
- l.17 $= \gamma^2 \frac{\partial \rho}{\partial t} + \gamma^2 \nabla \cdot (\rho \vec{u})$
- p.160 l.18 $T^{ab} = Au^a u^b + Bg^{ab}$ (7.8)
- l.22 $T^{ij} = \delta^{ij} P$ (7.9)
- p.161 l.3 $T^{ij} = B\eta^{ij}$
- l.6 $= Au^0 u^0 + B\eta^{00} =$
- l.9 $= (\rho + P)u^a u^b - P\eta^{ab}$ (7.10)
- l.11 $= (\rho + P)u^a u^b - Pg^{ab}$ (7.11)
- l.14 $= (\rho + P)u^a u^b + P\eta^{ab}$ (7.12)
- l.15 $= (\rho + P)u^a u^b + Pg^{ab}$
- p.162 l.5 $G_{\hat{a}\hat{b}} - \Lambda\eta_{\hat{a}\hat{b}} = 8\pi T_{\hat{a}\hat{b}}$ (7.14)
- l.13 Putting this together with (7.14) and (7.13), we have
- p.165 l.10 worldlines

Chapter 8

p.168, example 8-1: Alternative solution (*):

Calculate the Lie-derivative of g_{ab} (4.28) and using (4.18) you get:

$$\begin{aligned} L_X g_{ab} &= X^c \nabla_c g_{ab} + g_{cb} \nabla_a X^c + g_{ac} \nabla_b X^c \\ &= \nabla_a (g_{cb} X^c) + \nabla_b (g_{ac} X^c) \\ &= \nabla_a X_b + \nabla_b X_a \end{aligned}$$

If $L_X g_{ab} = 0$ this implies that $\nabla_a X_b + \nabla_b X_a = 0$, which is the Killing equation.

- p.171 l.6 $\nabla_\theta X_\theta = \partial_\theta X_\theta - \Gamma^c_{\theta\theta} X_c = \partial_\theta X_\theta - \Gamma^\theta_{\theta\theta} X_\theta - \Gamma^\phi_{\theta\theta} X_\phi$
- p.172 l.17 $\partial_\theta X_\phi = \partial_\theta [-\sin\theta \cos\theta \int f(\phi') d\phi' + g(\theta)]$
- p.174 l.8 $P(t) = \int p(t') dt'$
- l.10 $y(t) = e^{-P(t)} \int e^{P(t')} r(t') dt' + C e^{-P(t)}$
- l.14 $P(\theta) =$
- l.16 $e^{-P(\theta)} =$
- l.17 From this we deduce that $e^{P(\theta)} =$
- p.176 l.20 $X^\theta = g^{\theta\theta} X_\theta = \frac{1}{a^2} X_\theta$
- l.22 $X^\phi = g^{\phi\phi} X_\phi = \frac{1}{a^2} \frac{1}{\sin^2 \theta} X_\phi$
- p.177 l.6 $= \frac{1}{a^2} (A \cos \phi + B \sin \phi) \partial_\theta + \frac{1}{a^2} [C - \cot \theta (A \sin \phi - B \cos \phi)] \partial_\phi$
- l.7 $= A' \cos \phi \partial_\theta - A' \cot \theta \sin \phi \partial_\phi + B' \sin \phi \partial_\theta + B' \cot \theta \cos \phi \partial_\phi + C' \partial_\phi$
- l.8 $= +A' L_x + B' L_y + C' L_z$
- l.10 $L_x = +\cos \phi \frac{\partial}{\partial \theta} - \cot \theta \sin \phi \frac{\partial}{\partial \phi}$
- p.178 l.5 $\nabla_b T^{ab} = 0$

Chapter 9

- p.188 l.2 $= -r^2 \sin^2 \theta$

- p.190 l.2 $m \cdot \bar{m} = m^a \cdot \bar{m}_a =$
 l.23 $\bar{\delta} = \bar{m}^a \nabla_a$
- p.198 (*) l.22 $\Gamma^{vu}_{uu} = \frac{1}{2} g^{vd} (\partial_u g_{du} + \partial_u g_{ud} - \partial_d g_{uu}) = \frac{1}{2} g^{vu} (\partial_u g_{uu} + \partial_u g_{uu} - \partial_u g_{uu})$
 $= \frac{1}{2} g^{vu} (\partial_u g_{uu}) = \frac{1}{2} (1) \left(\frac{\partial H}{\partial u} \right) = + \frac{1}{2} \frac{\partial H}{\partial u}$ (9.29)
- p.199 l.8 Now $\nabla_b l_a =$ and so we can immediately...
 l.22 Having a look at the Christoffel symbols (9.29), we see ...

Chapter 10

- p.204 l.20 This tells us that the **metric** line element will not
- p.205 l.21 $= Cr^2 d\theta^2 + Cr^2 \sin^2 \theta d\phi^2$
 l.26 $= \sqrt{C} \left(\frac{r}{2c} \frac{dC}{dr} + 1 \right) dr$
- p.207 (*) l.19 $\Gamma^{\hat{t}}_{\hat{t}} = \Gamma^{\hat{t}}_{\hat{\theta}} = \Gamma^{\hat{t}}_{\hat{\phi}} = 0$ (10.16)
- p.208 (*) l.5 $\Gamma^{\hat{\theta}}_{\hat{t}} = \Gamma^{\hat{\theta}}_{\hat{\theta}} = \Gamma^{\hat{\theta}}_{\hat{\phi}} = 0$ (10.17)
- (*) l.9 $\Gamma^{\hat{\phi}}_{\hat{t}} = \Gamma^{\hat{\phi}}_{\hat{\phi}} = 0$ (10.18)
- p.210 l.12 $= -R^{\hat{r}}_{\hat{t}\hat{t}\hat{t}}$
 l.17 All together, the **independent** nonzero components of the
 l.19 $= -R^{\hat{t}}_{\hat{\phi}\hat{t}\hat{\phi}}$
- p.211 l.4 $R_{\hat{a}\hat{b}} = R^{\hat{c}}_{\hat{a}\hat{c}\hat{b}}$
- p.216 l.32 **l**agrangian
- p.217 (*) l.18 $= -\frac{2}{1 - \frac{2m}{r}} \ddot{r} + \frac{4m}{(r - 2m)^2} (\dot{r})^2$
- p.218 (*) l.2 $= \frac{2m}{r^2} (\dot{t})^2 + \left(1 - \frac{2m}{r}\right)^{-2} \left(\frac{2m}{r^2}\right) (\dot{r})^2 - 2r(\dot{\theta})^2 - 2r \sin^2 \theta (\dot{\phi})^2$
- p.220 l.19 $= -r^2 \sin^2 \theta \frac{d\phi}{d\tau} = -r^2 \frac{d\phi}{d\tau}$ (10.48)
- p.226 l.2 respect to ϕ r a second time
 l.19 ... and using $\mathcal{Y}r_0 = r \sin \phi$
 l.20 $\dots \frac{1}{A} = r \sin \phi = \mathcal{Y}r_0$
 l.21 the constant $1/A$ represents the
- p.227 l.5 $u_1^p = D \sin^2 \phi + E \cos^2 \phi$
 l.7 $u_1^p =$
 l.8 $u_1^p =$
 l.10 This line should be omitted: $u_p^t + u_p = -D \sin^2 \phi + 4D \sin^2 \phi = 3D \sin^2 \phi$
 l.11 $u_1^{p''} + u_1^p =$
 l.14 $u_1^{p''} + u_1^p = A^2 \sin^2 \phi$
 l.17 $u_1^{p''} + u_1^p = A^2 \sin^2 \phi$
- p.228 l.17 and setting **this** u equal to zero
- p.230 l.2 so we put ct in place of t and $\frac{Gm}{c^2 r}$ in place of $\frac{m}{r}$

$$\begin{aligned}
 \text{l.3} \quad ct &= \frac{dr}{\sqrt{1 - \frac{r_0^2}{r^2}}} \left(1 + \frac{2m}{r} \frac{G}{c^2} - \frac{mr_0^2}{r^3} \frac{G}{c^2} \right) \\
 \text{l.7} \quad ct &= \sqrt{r_p^2 - r_0^2} - \sqrt{r_e^2 - r_0^2} + 2m \frac{G}{c^2} \ln \frac{(r_p + \sqrt{r_p^2 - r_0^2})(r_e + \sqrt{r_e^2 - r_0^2})}{r_0^2} \\
 &\quad - m \frac{G}{c^2} \left[\frac{\sqrt{r_p^2 - r_0^2}}{r_p} - \frac{\sqrt{r_e^2 - r_0^2}}{r_e} \right]
 \end{aligned}$$

l.20 Using the variational method described in Example 4-10, the nonzero Christoffel symbols for the **general** Schwarzschild metric are

p.231 (*) Q2 (b) $R_{rt} = \frac{2}{r} \left(\frac{d\lambda}{dt} \right)$

(*) Q3 (b) $\Gamma^{\hat{r}}_{\hat{t}\hat{t}} = \frac{3m - \Lambda r^3}{r^{3/2} \sqrt{9r - 18m - 3\Lambda r^3}}$

p.236 l.27 **Eddington-Finkelstein Coordinates** The Radial Null Geodesic

p.238 (*) l.10 $t - t_0 = -\frac{2}{3\sqrt{2m}} \left(r^{3/2} - r_0^{3/2} + 6m\sqrt{r} - 6m\sqrt{r_0} \right) + 2m \dots$

p.239 (*) l.2 $r - 2m = (r_0 - 2m)e^{-(t-t_0)/2m}$

Chapter 11

p.245 l.26 $ds^2 = \left(1 - \frac{2m}{\Sigma} \right) dt^2 + \frac{4amr \sin^2 \theta}{\Sigma} dt d\theta - \dots$ (11.9)

p.246 (*) l.12 $g^{\phi\phi} = -\frac{(\Delta - a^2 \sin^2 \theta)}{\Sigma \Delta \sin^2 \theta}$ (11.13)

p.247 l.4 In this case, $\theta = \frac{\pi}{2}$, which
 l.5 means that **sin θ = 1** and

p.251 l.3 $R^2 = r^2 + a^2 + \frac{2m^3}{r}$ (11.19)

Chapter 12

p.258 l.11 Suppose that **S** represents the spacelike...

p.261 l.4 $e^{2f} = \frac{1}{C' - kr^2}$ (12.4)

l.5 To find the constant $C' = -2C$, we can use the ...

l.13 $= 1 - C' + Kr^2 - Kr^2 = 1 - C' + 2Kr^2$

l.15 $1 - C' + 2Kr^2 = 2Kr^2$ or $1 - C' = 0 \Rightarrow C' = 1$

p.268 l.13 The detail were worked out in Example 7-3

P.273 l.20 $ds^2 = dt^2 - e^{2\sqrt{\frac{\Lambda}{3}}t} [dr^2 + r^2 d\theta^2 - r^2 \sin^2 \theta d\phi^2]$

p.274 Fig. 12-5 The vertical coordinate is **a** not **R**

l.1 The de Sitter solution represents a universe without matter **and without radiation**.

p.275 Fig. 12-6 The vertical coordinate is **a** not **R**

l.10 $\frac{da}{dt} = 2C \sin \tau \cos \tau \frac{d\tau}{dt}$

p.276 Fig. 12-7 The vertical coordinate is **a** not **R**

l.3 Integrating the **left** side becomes

p.277 (*) Q2 $re^{2\nu(r)} = A + Br + \frac{1}{3}k^2\Lambda r^3$

(*) Q3 $B = k^2$

(*) Q4 $dl^2 = \frac{dr^2}{1 - \frac{2m}{r} + \frac{1}{3}\Lambda r^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$

Chapter 13

p.280 l.25 $\Gamma^a_{bc} = \frac{1}{2}g^{ad} \left(\frac{\partial g_{bd}}{\partial x^c} + \frac{\partial g_{cd}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^d} \right)$

p.282 l.6 $\Gamma^a_{bc} = \frac{1}{2}g^{ad} \left(\frac{\partial g_{bd}}{\partial x^c} + \frac{\partial g_{cd}}{\partial x^b} - \frac{\partial g_{bc}}{\partial x^d} \right)$

l.7 $= \frac{\varepsilon}{2}g^{ad} \left(\frac{\partial h_{bd}}{\partial x^c} + \frac{\partial h_{cd}}{\partial x^b} - \frac{\partial h_{bc}}{\partial x^d} \right)$

l.8 $= \frac{\varepsilon}{2}(\eta^{ad} - \varepsilon h^{ad}) \left(\frac{\partial h_{bd}}{\partial x^c} + \frac{\partial h_{cd}}{\partial x^b} - \frac{\partial h_{bc}}{\partial x^d} \right)$

l.9 $= \frac{1}{2}(\varepsilon\eta^{ad} - \varepsilon^2 h^{ad}) \left(\frac{\partial h_{bd}}{\partial x^c} + \frac{\partial h_{cd}}{\partial x^b} - \frac{\partial h_{bc}}{\partial x^d} \right)$

l.12 $= \frac{1}{2}\varepsilon\eta^{ad} \left(\frac{\partial h_{bd}}{\partial x^c} + \frac{\partial h_{cd}}{\partial x^b} - \frac{\partial h_{bc}}{\partial x^d} \right)$ (13.3)

p.283 l.3 $= \partial_c \left[\frac{1}{2}\varepsilon\eta^{ae} \left(\frac{\partial h_{be}}{\partial x^d} + \frac{\partial h_{de}}{\partial x^b} - \frac{\partial h_{bd}}{\partial x^e} \right) \right]$

l.4 $- \partial_d \left[\frac{1}{2}\varepsilon\eta^{af} \left(\frac{\partial h_{bf}}{\partial x^c} + \frac{\partial h_{cf}}{\partial x^b} - \frac{\partial h_{bc}}{\partial x^f} \right) \right]$

l.5 $= \frac{1}{2}\varepsilon \left(\eta^{ae} \frac{\partial^2 h_{be}}{\partial x^c \partial x^d} + \eta^{ae} \frac{\partial^2 h_{de}}{\partial x^c \partial x^b} - \eta^{ae} \frac{\partial^2 h_{bd}}{\partial x^c \partial x^e} \right)$

l.6 $- \frac{1}{2}\varepsilon \left(\eta^{af} \frac{\partial^2 h_{bf}}{\partial x^d \partial x^c} + \eta^{af} \frac{\partial^2 h_{cf}}{\partial x^d \partial x^b} - \eta^{af} \frac{\partial^2 h_{bc}}{\partial x^d \partial x^f} \right)$

l.7 Let's relabel it as f

l.12 The **first** term will cancel

l.13 $= \frac{1}{2}\varepsilon\eta^{ae} \left(-\frac{\partial^2 h_{bd}}{\partial x^c \partial x^f} + \frac{\partial^2 h_{df}}{\partial x^c \partial x^b} + \frac{\partial^2 h_{bc}}{\partial x^d \partial x^f} - \frac{\partial^2 h_{cf}}{\partial x^d \partial x^b} \right)$ (13.4)

l.15 $R_{ab} = \partial_c \Gamma^c_{ab} - \partial_b \Gamma^c_{ac}$

p.287 l.2 $\psi'^a_{b,a} = \psi^a_{b,a} - \square\phi_b$

l.4 $W\phi_a = \psi^b_{a,b}$

$\psi'^a_{b,a}$

l.6 $\frac{1}{2}\varepsilon W\psi'_{ab} =$

l.8 $W\psi'_{ab} = 0$

l.14 $W\psi'_{ab} = W \left(h'_{ab} - \frac{1}{2}\eta_{ab}h' \right) = Wh'_{ab} - \frac{1}{2}W\eta_{ab}h' = 0$

l.16 $0 = W\psi'_{ab} = \eta^{ab}W\psi'_{ab} = W(\eta^{ab}\psi'_{ab}) = W(\psi'^b_b) = W\psi' = -Wh'$

l.19 $Wh'_{ab} = 0$

p.288 l.13 tensor can be shown to be a function of h_{xx}, h_{xy}, h_{yx} and h_{yy} alone, in the **Einstein gauge**.

p.289 l.1 derivatives with respect to y and zx vanish

p.293 (*) l.11 $ds^2 = d^2t - d^2x - d^2y - d^2z + 2\varepsilon h_{xy} dx dy$ (13.19)

p.294 (*) l.2 $ds^2 = dt^2 - (1 - \varepsilon h_{xy})dx'^2 - (1 + \varepsilon h_{xy})dy'^2 - dz^2$ (13.20)

p.299 l.5 **{see Quiz}**

- p.300 l.11 $-2 \frac{a'}{a} X du dX$
 l.12 $= b^2 \left(-\frac{b'}{b} Y du + \frac{1}{b} dY \right)^2 = \dots - 2 \frac{b'}{b} Y du dY \dots$
- p.303 l.4 The ~~curvature tensor~~ Weyl scalar
 p.304 l.11 $ds^2 = \delta(u)(X^2 - Y^2)du^2 + 2dudr - dX^2 - dY^2$ (13.42)
 l.13 ~~(see Problem 7)~~
- p.305 l.15 (13.45)
 p.306 l.13 $= -\frac{1}{2} \partial_\nu g_{xx}$
 l.14 $= -\frac{1}{2} \frac{\partial}{\partial \nu} (-[1 - v\theta(v)]^2)$
 l.15 $= (1 - v\theta(v)) \frac{\partial}{\partial \nu} (-v\theta(v))$
 l.16 $= (1 - v\theta(v)) \left(-\theta(v) - v \frac{d\theta}{dv} \right)$
 l.17 $= -(1 - v\theta(v)) (\theta(v) - v \delta(v))$
 l.19 $\int_{-\infty}^{\infty} f(v) \delta(v) dv = f(0)$
- p.307 l.7 since n_u is ~~not~~ 1
 l.8 $n_a = (1, 0, 0, 0)$
- p.308 l.3 $m^a = \frac{1}{\sqrt{2}} \left(0, 0, -\frac{1}{(1 - v\theta(v))}, -i \frac{1}{(1 + v\theta(v))} \right)$
 l.4 $\bar{m}^a = \frac{1}{\sqrt{2}} \left(0, 0, -\frac{1}{(1 - v\theta(v))}, i \frac{1}{(1 + v\theta(v))} \right)$
- p.309 l.2 $\bar{m}^x \bar{m}^x = \left(-\frac{1}{\sqrt{2}(1 - v\theta(v))} \right) \left(-\frac{1}{\sqrt{2}(1 - v\theta(v))} \right)$
- p.311 fig 13-11 $u^2 + v^2 = 1$
- p.315 (*) l.6 $m_a = \frac{1}{\sqrt{2}} (0, 0, \cos av, +i \cosh av)$ (13.56)
- p.317 l.1 $\sigma = \frac{a}{2} (\tan av + \tanh av)$
 l.6 $\Psi_0 = +D\sigma - 2\sigma\rho$ (13.57)
 l.11 $= \frac{\partial}{\partial \nu} \left[\frac{a}{2} (\tan av + \tanh av) \right]$
 (*) l.12 $= \frac{a^2}{2} (1 + \tan^2 av + 1 - \tanh^2 av)$
 l.14 $\Psi_0 = +D\sigma - 2\sigma\rho$
 (*) l.15-18 $= \frac{a^2}{2} (1 + \tan^2 av + 1 - \tanh^2 av) - \dots = a^2$
- p.318 l.7 *Nariai*
 l.8 $R_{ab} = -\Lambda g_{ab}$
 (*) l.11 if $R_{ab} = -g_{ab} \Lambda$ Ω should rightfully be $\Omega = 1 + \frac{\Lambda}{4} (x^2 + y^2)$
 l.16 $g^{\nu\nu} = +\Lambda v^2$
 (*) l.18 $\Gamma^{\nu}_{uv} = -\Gamma^u_{uu} = -\Lambda v$; $\Gamma^{\nu}_{uu} = \Lambda^2 v^3$; $\Gamma^x_{xx} = \Gamma^y_{yx} = -\Gamma^x_{yy} = -\frac{\Lambda x}{\Omega}$; $\Gamma^x_{xy} = -\Gamma^y_{xx}$
 $= \Gamma^y_{yy} = -\frac{\Lambda y}{\Omega}$

- l.23 $m_a = \frac{1}{\sqrt{2}} \left(0, 0, -\frac{1}{\Omega}, -i\frac{1}{\Omega} \right); \bar{m}_a = \frac{1}{\sqrt{2}} \left(0, 0, -\frac{1}{\Omega}, \frac{i}{\Omega} \right)$
- l.24-25 These lines should be omitted
- p.319 l.3 $n^a = \left(1, +\frac{1}{2}\Lambda v^2, 0, 0 \right)$
- l.4 $m^a = \frac{1}{\sqrt{2}} (0, 0, \Omega, i\Omega); \bar{m}^a = \frac{1}{\sqrt{2}} (0, 0, \Omega, -i\Omega)$
- l.11 $-\Lambda v$
- l.12 $l_{a;b} n^a n^b = l_{u;u} n^u n^u = -\Lambda v$
- l.19 $m_x = -\frac{i}{\sqrt{2}\Omega}$
- l.22 $\gamma = \frac{1}{2} l_{u;u} n^u n^u = -\frac{1}{2} \Lambda v$
- p.320 (*) l.1 An exercise shows that **all** remaining nonzero spin coefficients **vanish** are $\alpha =$
 $-\frac{\Lambda}{\sqrt{2}}(x - iy); \beta = \frac{\Lambda}{\sqrt{2}}(x + iy)$
- l.5 $= \rho\mu - \sigma\lambda + \alpha\bar{\alpha} + \beta\bar{\beta} - 2\alpha\beta + \gamma(\rho - \bar{\rho}) + \epsilon(\mu - \bar{\mu}) - \Psi_2 + \Lambda_{NP} + \Phi_{11} \quad (13.59)$
- l.8 $= \partial_v \left(-\frac{1}{2} \Lambda v \right) = -\frac{1}{2} \Lambda$
- l.10 $= -\frac{1}{2} \Lambda \quad (13.61)$
- l.18 $-\frac{1}{2} \Lambda =$
- l.19 $= -\frac{1}{12} \Lambda \quad (13.64)$
- p.321 l.2 $\Psi_2 = -\frac{1}{6} \Lambda$ and $\Phi_{11} = -\frac{1}{4} \Lambda \quad (13.65)$
- p.322 Q1 l.2 used in Example 13-2
- Q2 l.13 $= -2\sqrt{2}\mu \left(\frac{x}{x^2 + y^2} + i \frac{y}{x^2 + y^2} \right)$
- Q3 l.15 used in Example 13-1
- l.18 $\Phi_{22} = \delta v - \Delta\mu - \mu^2 - \lambda\bar{\lambda} - \mu(\gamma + \bar{\gamma}) + \bar{v}\pi - v(\tau - 3\beta - \bar{\alpha}) \quad (9.24)$

Final Exam

- p.323 E1 (*) l.2 $\Delta s^2 \dots E_1 = (-1, 3, 2, 4)$
- (*) l.3 $E_2 = (4, 0, -1, 1)$
- p.324 E6 (*) l.25 $\Gamma_{aba} = \Gamma_{aab} = \frac{1}{2} \frac{\partial g_{aa}}{\partial x^b}$
- E7 (*) The nonzero Christoffel symbols are: $\Gamma^u_{uu} = -\Gamma^u_{vv} = \Gamma^v_{uv} = \frac{u}{(u^2+v^2)}$; $\Gamma^v_{vv} = -\Gamma^v_{uu} = \Gamma^u_{vu} = \frac{v}{(u^2+v^2)}$; $\Gamma^u_{\theta\theta} = -\frac{uv^2}{(u^2+v^2)}$; $\Gamma^v_{\theta\theta} = -\frac{u^2v}{(u^2+v^2)}$; $\Gamma^\theta_{u\theta} = \frac{1}{u}$; $\Gamma^\theta_{v\theta} = \frac{1}{v}$
- p.325 E9 (*) l.11 $\Gamma^\psi_{\theta\theta} = -\cosh \psi \sinh \psi, \Gamma^\psi_{\phi\phi} = -\cosh \psi \sinh \psi \sin^2 \theta$
- E10 (*) The non-zero Ricci rotation coefficients are: $\Gamma^{\hat{\psi}}_{\hat{\theta}\hat{\theta}} = -\Gamma^{\hat{\theta}}_{\hat{\psi}\hat{\theta}} = -\coth \psi, \Gamma^{\hat{\psi}}_{\hat{\phi}\hat{\phi}} = -\Gamma^{\hat{\phi}}_{\hat{\theta}\hat{\phi}} = -\frac{\cot \theta}{\sinh^2 \psi}$
- p.326 E12 (*) l.2 $\Psi_2 = \frac{r^2 - 4rm + 3e^2}{2r^4}$

$$\begin{aligned}
 \text{E14 (*)} \quad \text{l.16} \quad ds^2 &= \frac{1}{2\omega^2} \left((dt + e^x dz)^2 - dx^2 - dy^2 - \frac{1}{2} e^{2x} dz^2 \right) \\
 \text{(*)} \quad \text{l.18} \quad T_{ab} &= \frac{\rho}{2\omega^2} \begin{pmatrix} 1 & 0 & 0 & e^x \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ e^x & 0 & 0 & e^{2x} \end{pmatrix}
 \end{aligned}$$

Quiz and Exam Solutions

p.330 l.3 (*) 7.b
 p.331 l.7 (*) 5.c
 l.9 (*) 13.a

Bibliografi

McMahon, D. (2006). *Relativity Demystified*. McGraw-Hill Companies, Inc.

¹ http://en.wikipedia.org/wiki/G%C3%B6del_metric